

Binary Evolution and its Applications

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Yunnan Observatories of the CAS

Collaborators:

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R. Webbink, Ph. Podsiadlowski, A. Lynas-Gray, L. Yungelson, M. Gilfanov, T.E. Wood, G. Stephan,

22-Mar-2018

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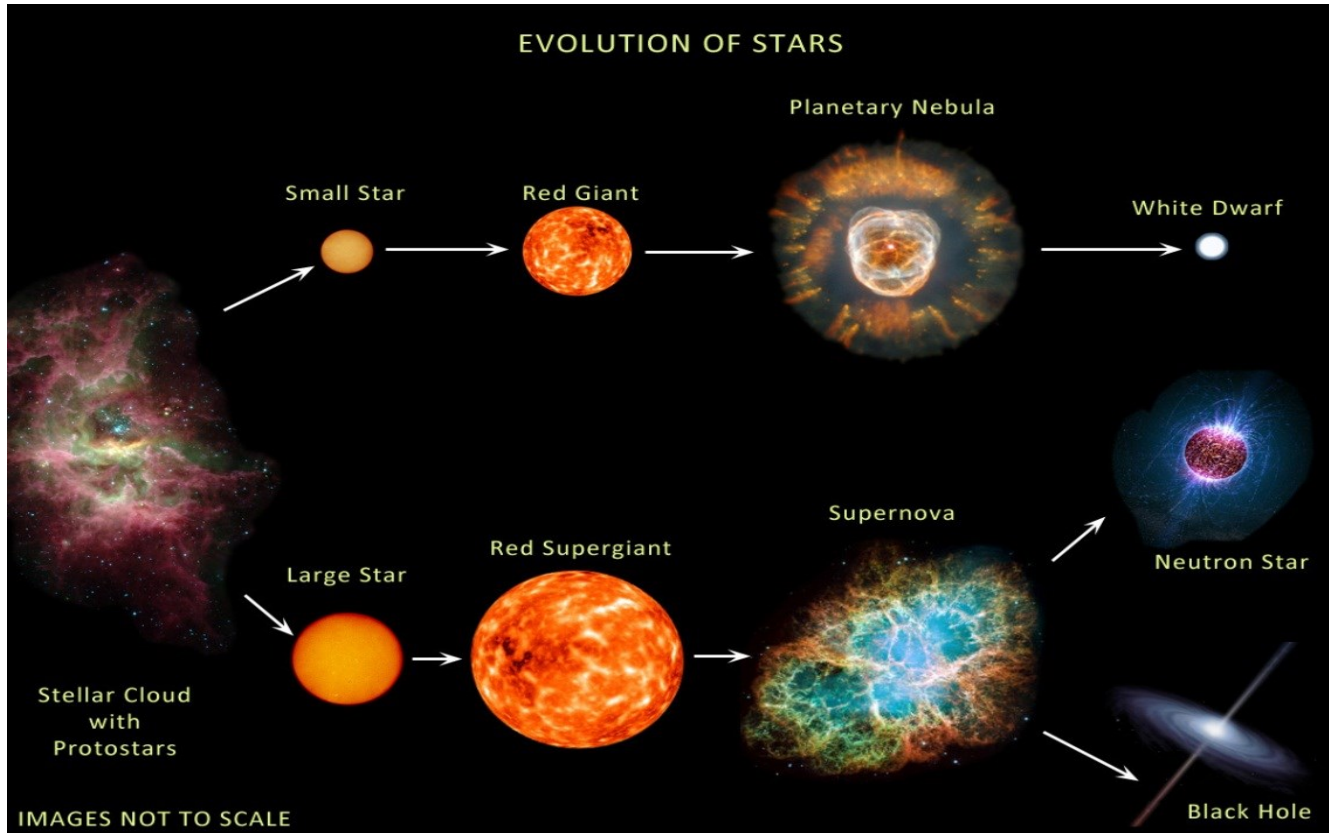
Outline

- Binary evolution
- Binary-related objects
- Evolutionary population synthesis
- Future perspectives

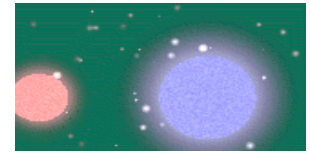
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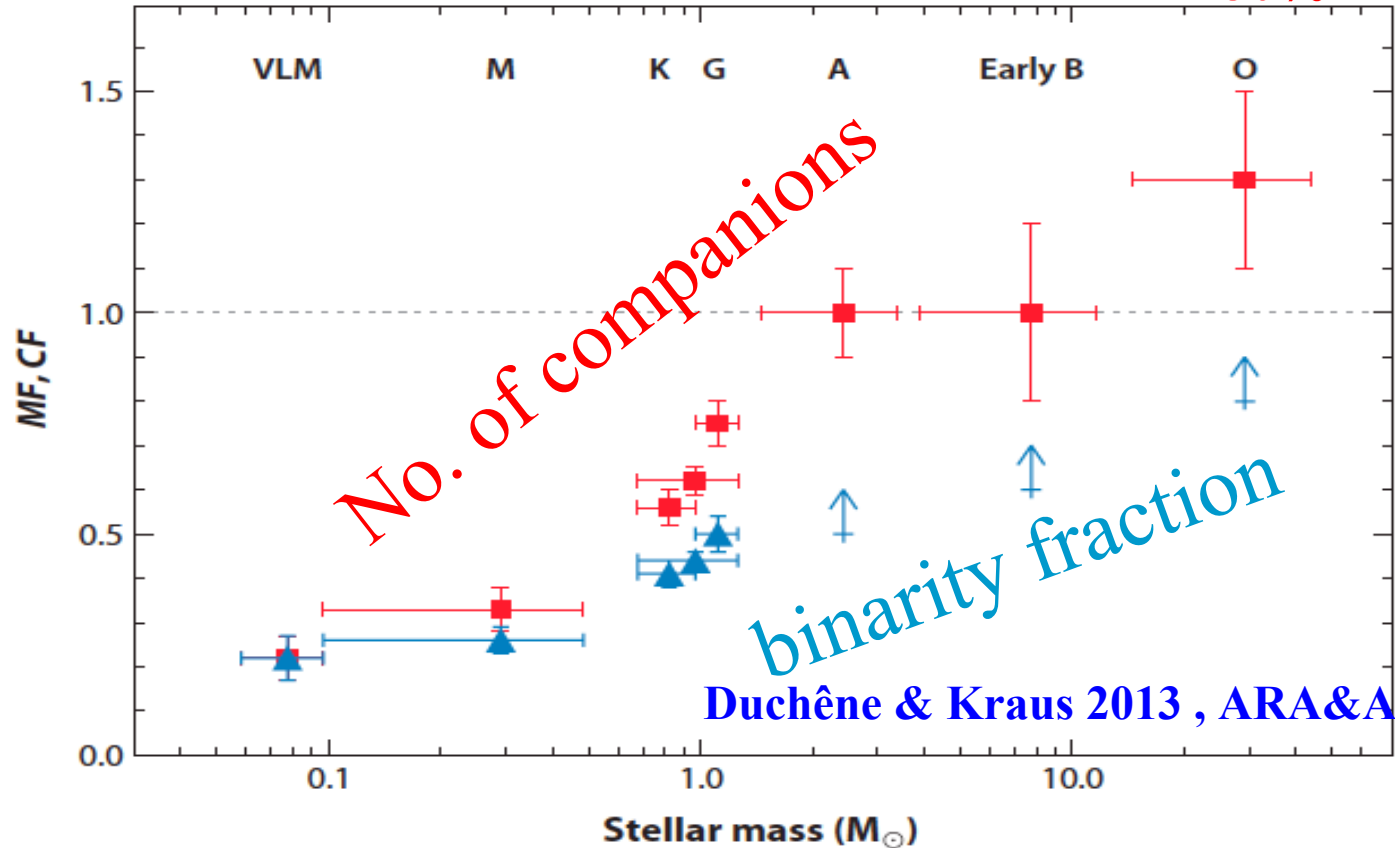
Stellar evolution



Half of the stars are in binaries



50% in binaries

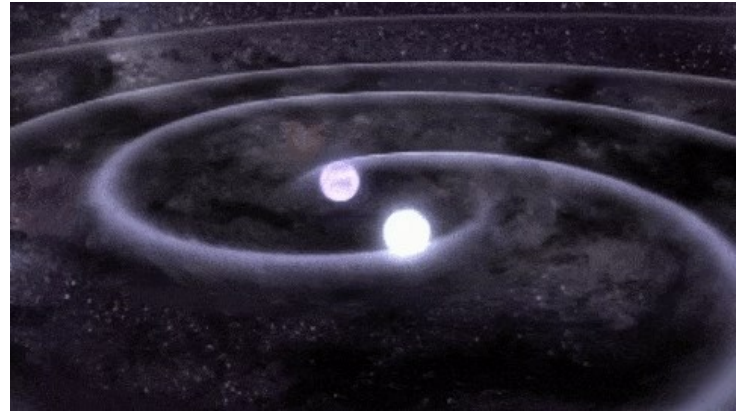
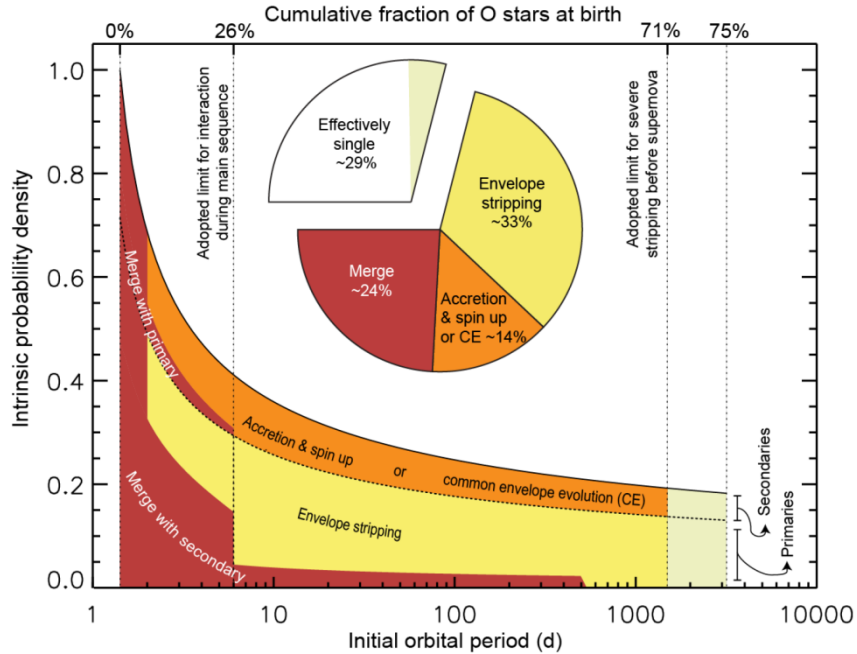


No. of companions

binary fraction

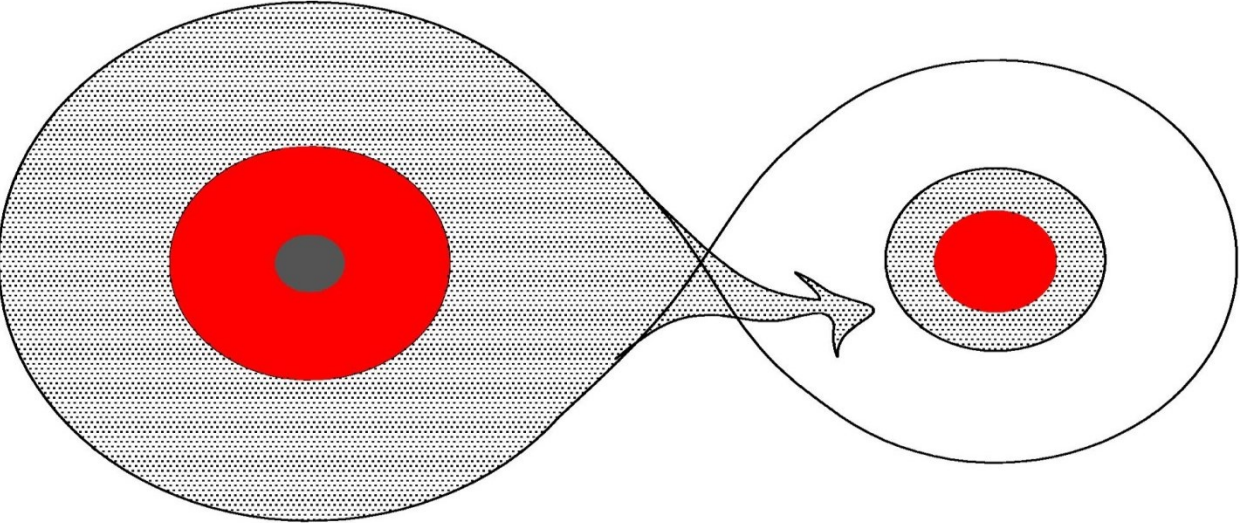
Duchêne & Kraus 2013, ARA&A

Binary interaction dominates the evolution of massive stars

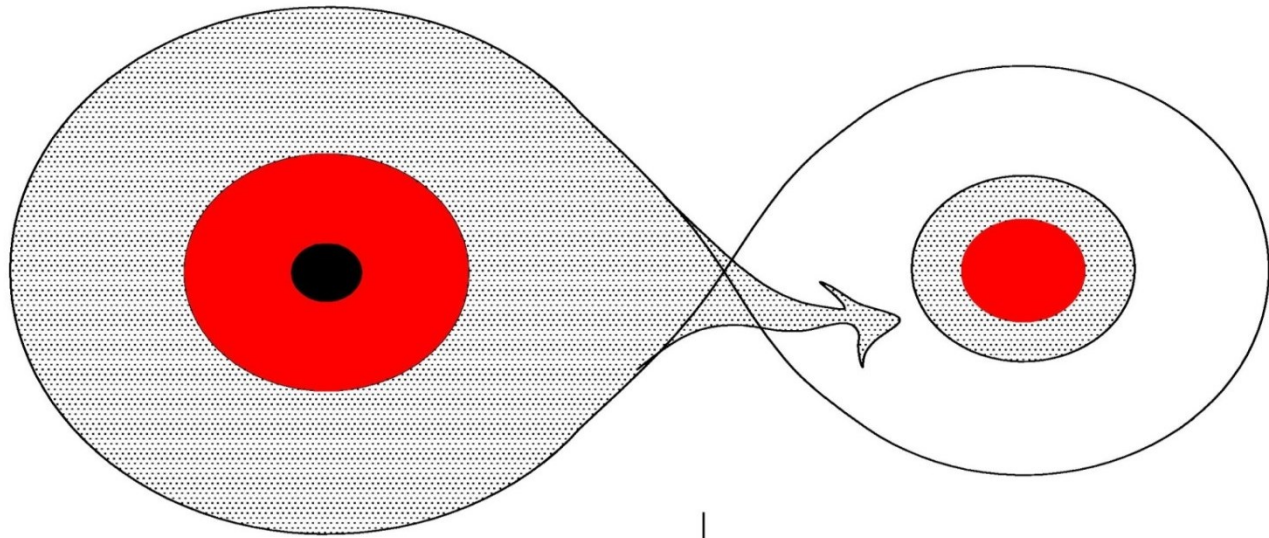


Sana et al 2012 Sci

Binary Evolution



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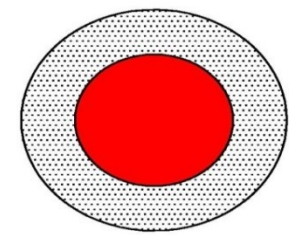


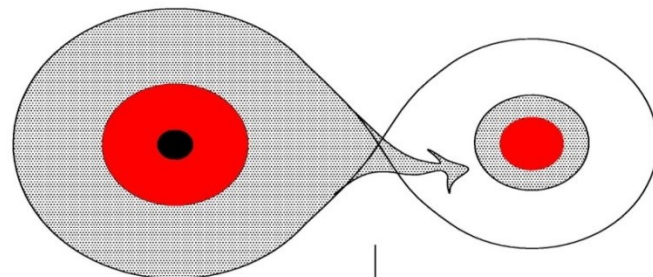
Stable RLOF



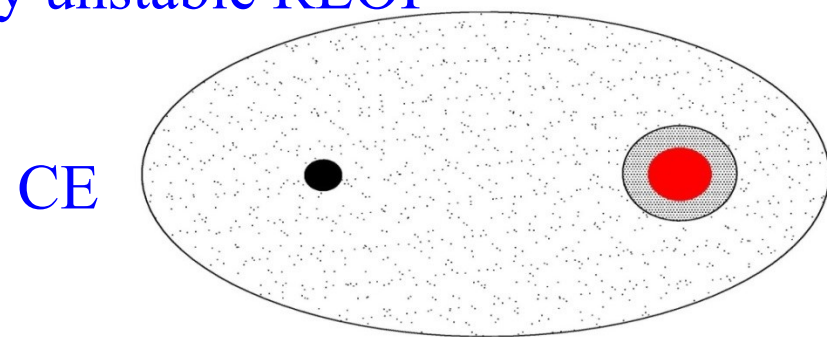
WD

MS



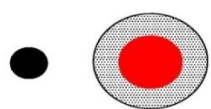


Dynamically unstable RLOF



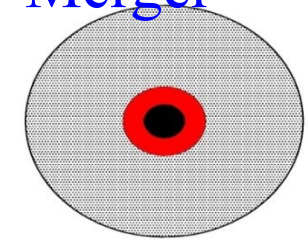
CE

Ejection

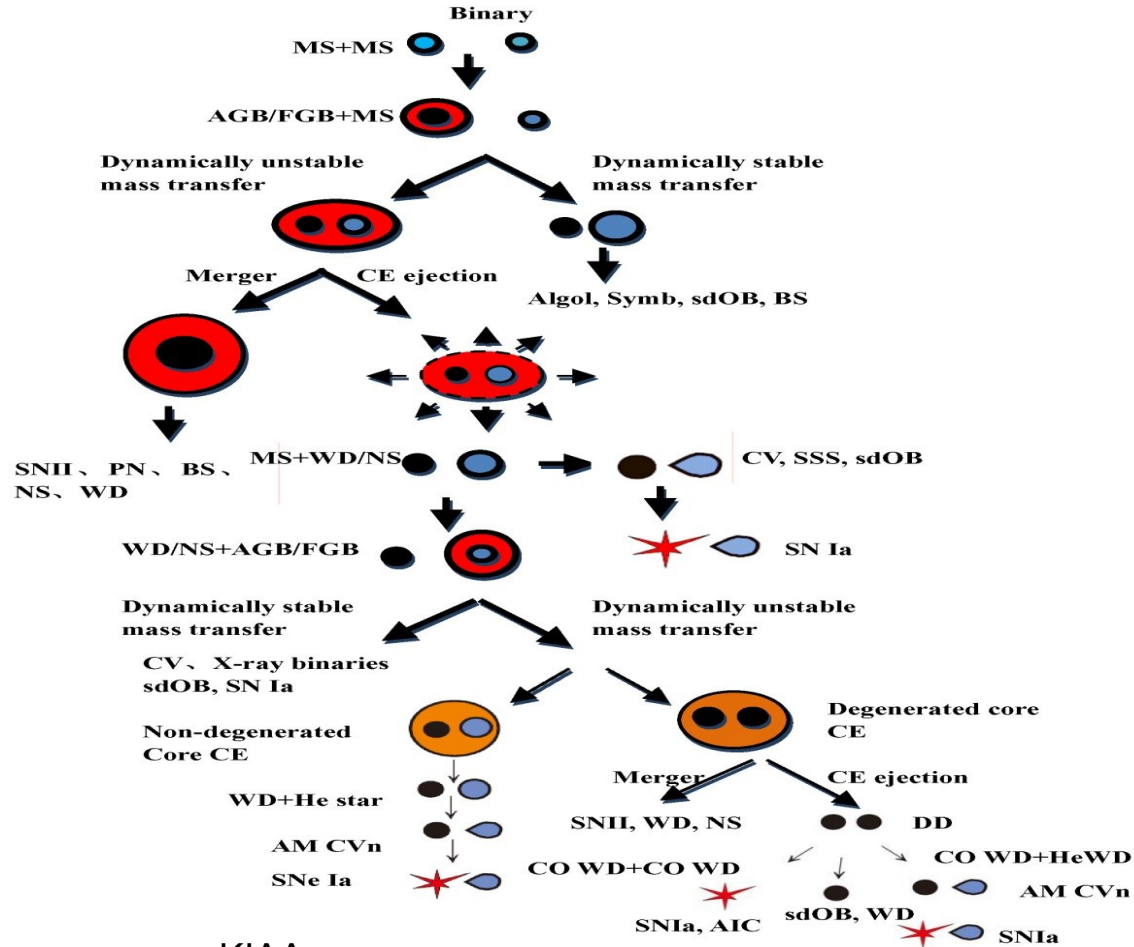


WD MS

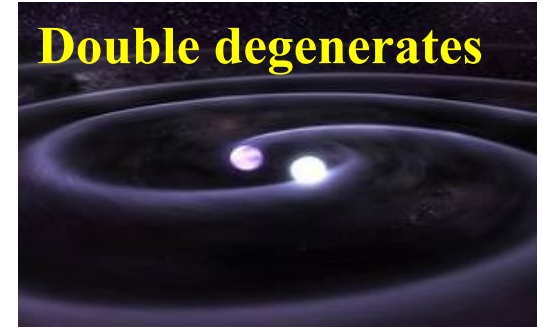
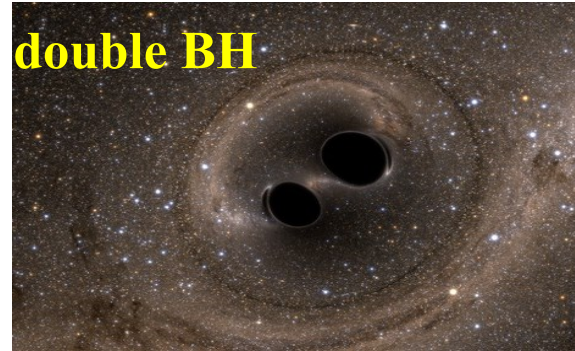
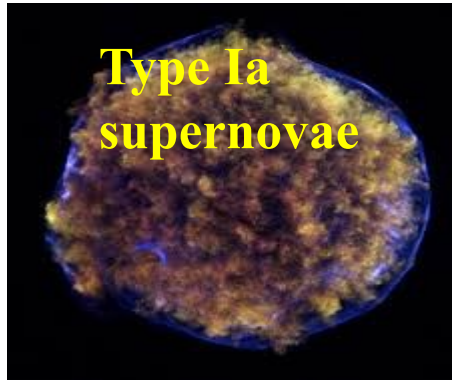
Merger



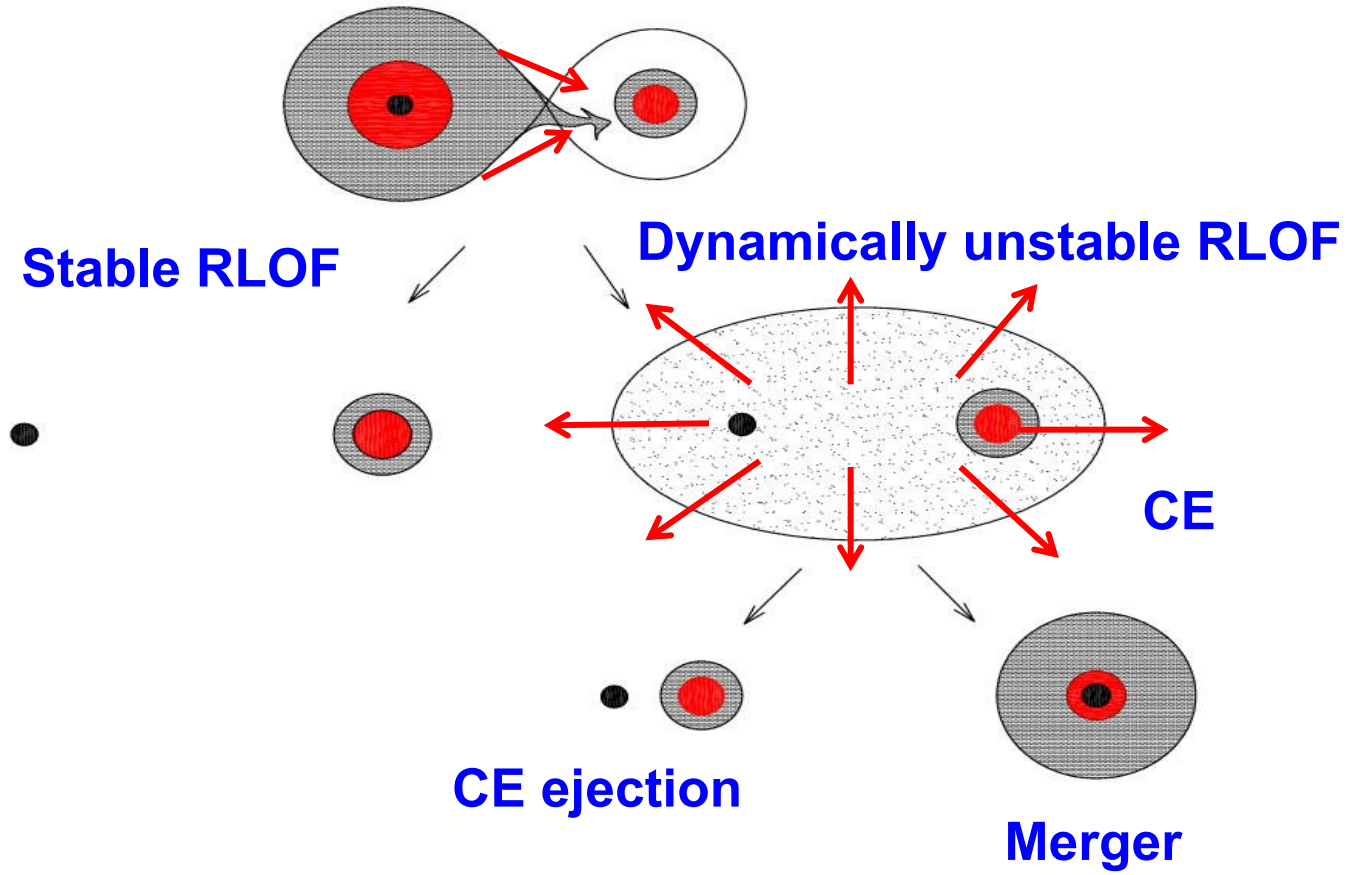
Binary Evolution Flow Chart



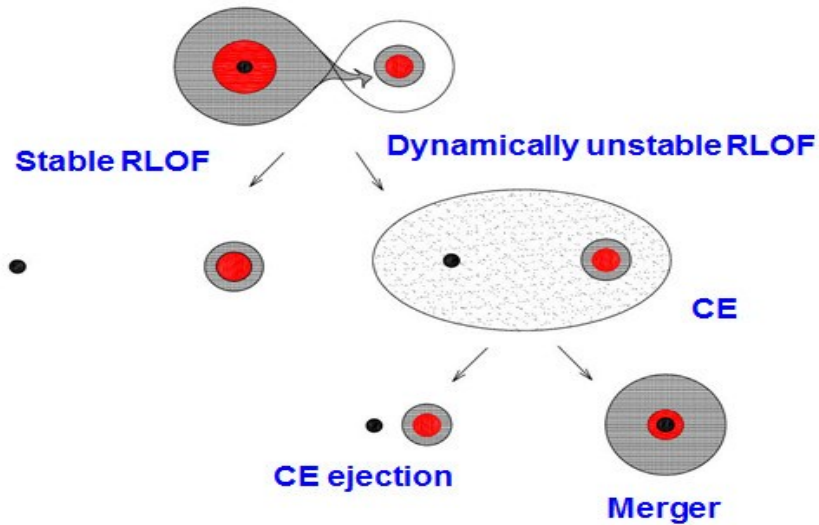
Binary evolution produces exotic objects



Basic problems:



Basic problems:



mass transfer stability?



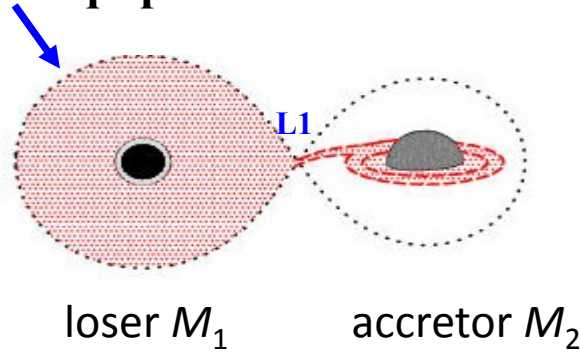
CE evolution?

RLOF Stability criterion models

- Polytropic model
(Hjellming+ 1987, Webbink 1988)
- detailed binary evolution calculation model (Pavlovskii+ 2015, 2017)
- Adiabatic mass loss model
(Ge+ 2010, 2015)

Mass transfer in Binaries

Critical equipotential surface



Roche lobe radius:

$$\frac{R_{L1}}{a} = \frac{0.49q^{2/3}}{0.6q^{2/3} + \ln(1 + q^{1/3})}$$

$$q = M_1/M_2$$

Eggleton, 1983, ApJ

Response of the loser

$$\zeta_{th} = \left. \frac{\partial \ln R_1}{\partial \ln M_1} \right|_{th}$$

$$\zeta_{ad} = \left. \frac{\partial \ln R_1}{\partial \ln M_1} \right|_{ad}$$

VS

Response of the Roche lobe

$$\zeta_L = \left. \frac{\partial \ln R_L}{\partial \ln M_1} \right|_{RLOF}$$

$$(\zeta_{ad}, \zeta_{th}) > \zeta_L$$

Nuclear Timescale MT (driven by nuclear burning)

$$\zeta_{ad} > \zeta_L > \zeta_{th}$$

Thermal timescale MT (driven by recovery of thermal equilibrium)

$$\zeta_L > \zeta_{ad}$$

Dynamical timescale MT (constrained by hydrostatic equilibrium only)

➤ **Stellar Structure Equations**

$$\frac{d \ln P}{dm} = -\frac{Gm}{4\pi r^4}$$

hydrostatic equilibrium

$$\frac{dr^2}{dm} = \frac{1}{2\pi r \rho'}$$

Mass continuity

$$\frac{d \ln T}{dm} = \frac{d \ln P}{dm} \nabla$$

$$\frac{dL}{dm} = \epsilon - \epsilon_\nu - C_P T \left(\frac{d \ln T}{dt} - \nabla_a \frac{d \ln P}{dt} \right)$$

Energy equations

$$\zeta_L > \zeta_{ad}$$

Dynamical timescale MT

➤ **Equations in the adiabatic mass loss model**

$$\frac{d \ln P}{dm} = -\frac{Gm}{4\pi r^4 P'}$$

$$\frac{dr^2}{dm} = \frac{1}{2\pi r \rho'}$$

Adiabatic assumption

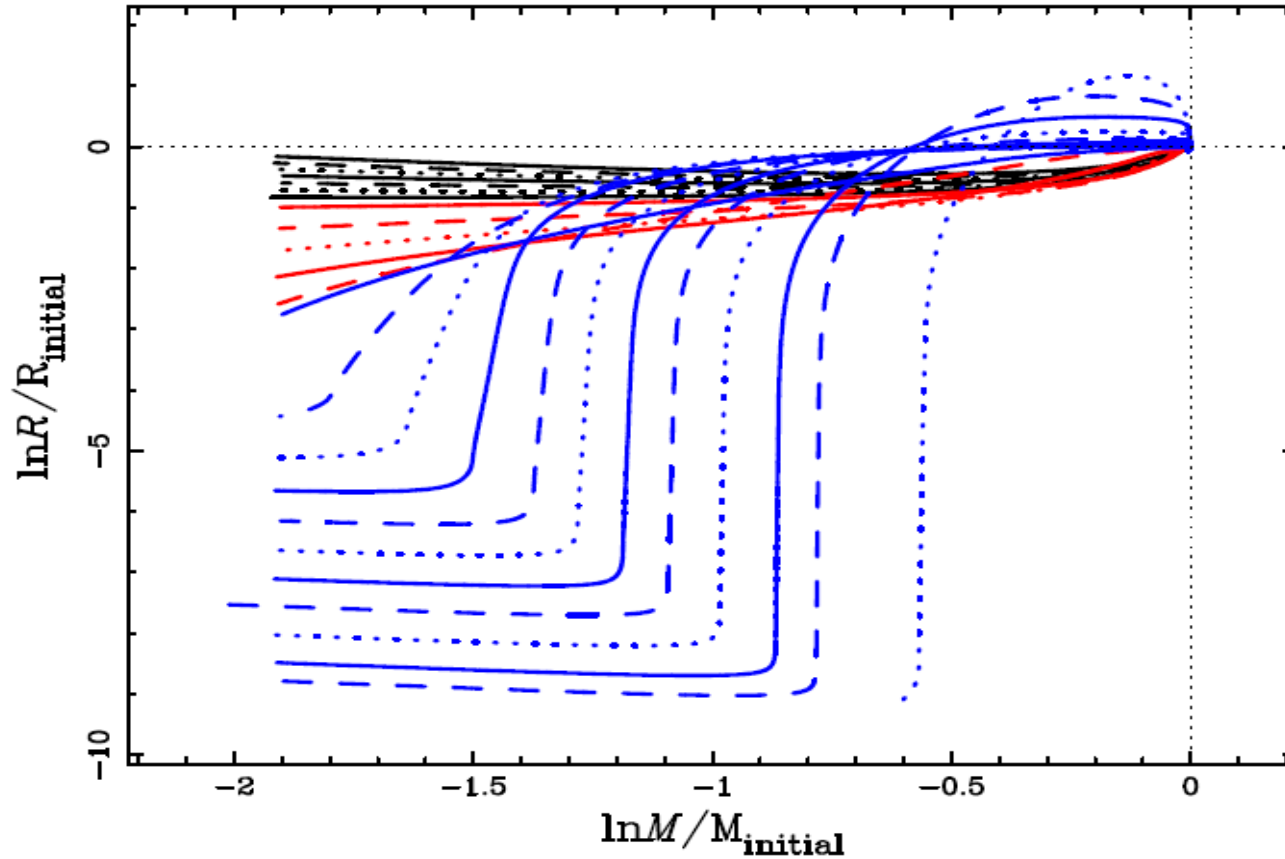
$$s' \equiv ds/dt = s(m) - s_0(m) = 0,$$

$$X' \equiv dX/dt = X(m) - X_0(m) = 0.$$

$$\zeta_{ad} = \left. \frac{\partial \ln R_1}{\partial \ln M_1} \right|_{ad}$$

Ge+, 2010

Response of a $1.00M_{\odot}$ star to adiabatic mass loss



q_c for FGB or AGB donors

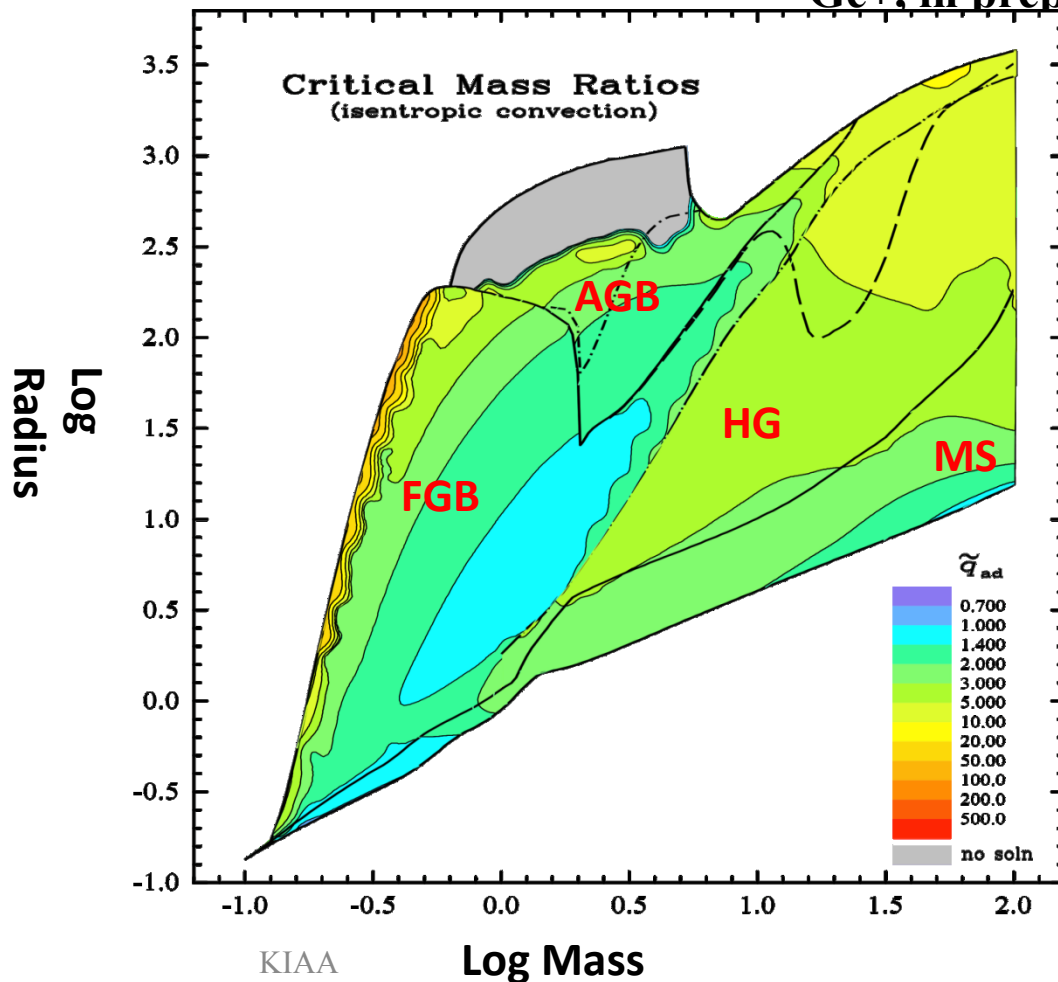
Ge+, 2015

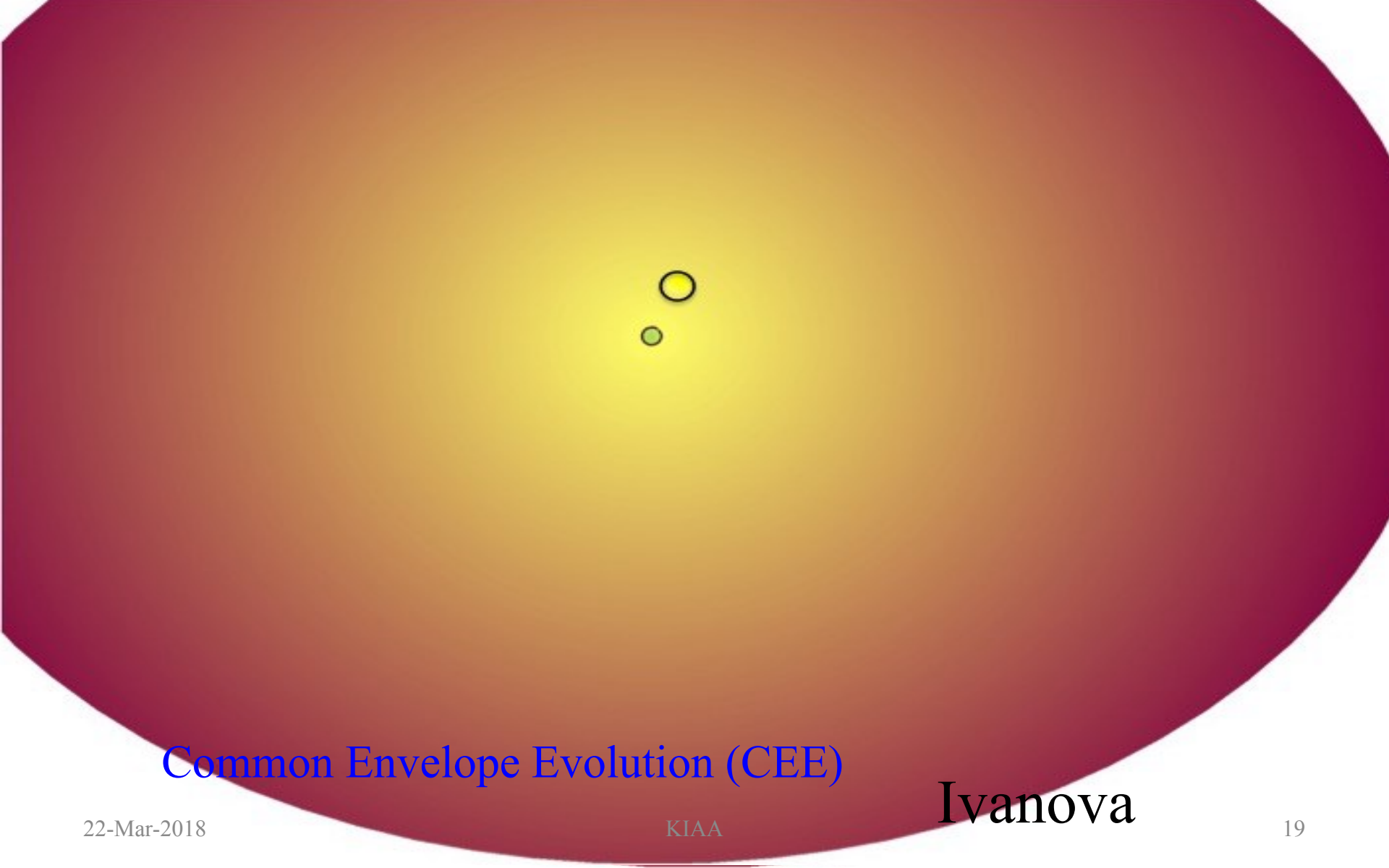
Ge+, in prep

Before: $q_c \sim 1$
CE only

Now: from 1.4-5
Both CE and RL

Non-conservative



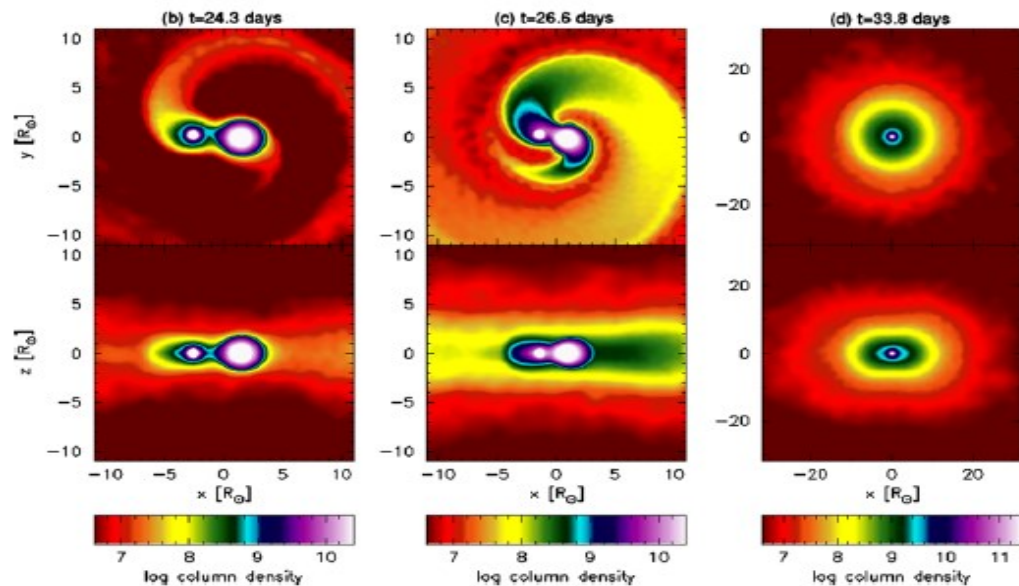


Common Envelope Evolution (CEE)

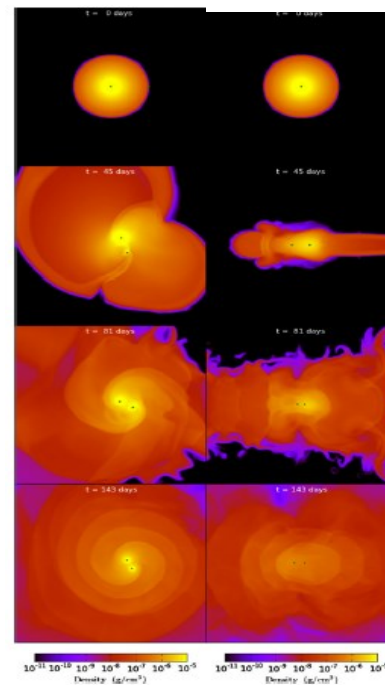
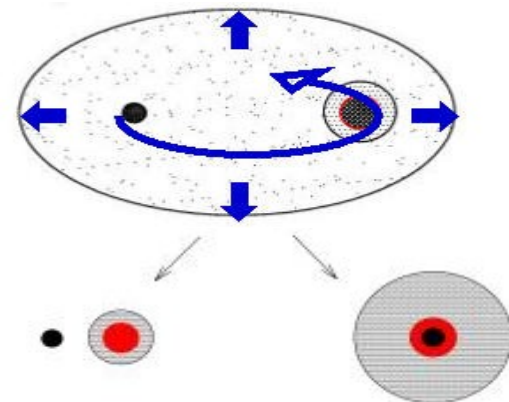
Ivanova

Common Envelope Evolution

SPH simulations are difficult !



merger



ejection

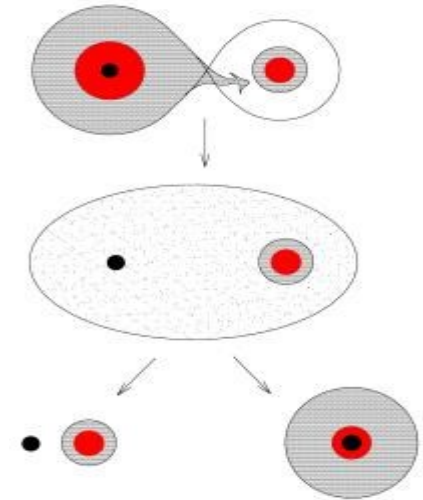
Common Envelope Evolution

Parameterized !

$$\Delta E_{\text{orb}} \simeq \frac{GM_c M_2}{2a_f} - \frac{G(M_c + M_e)M_2}{2a_i}$$

$$\alpha_{\text{CE}} \Delta E_{\text{orb}} \geq E_{\text{bind}} = E_{\text{gr}} - \alpha_{\text{th}} E_{\text{th}}$$

E_{gr} and E_{th} are integrated for the envelope

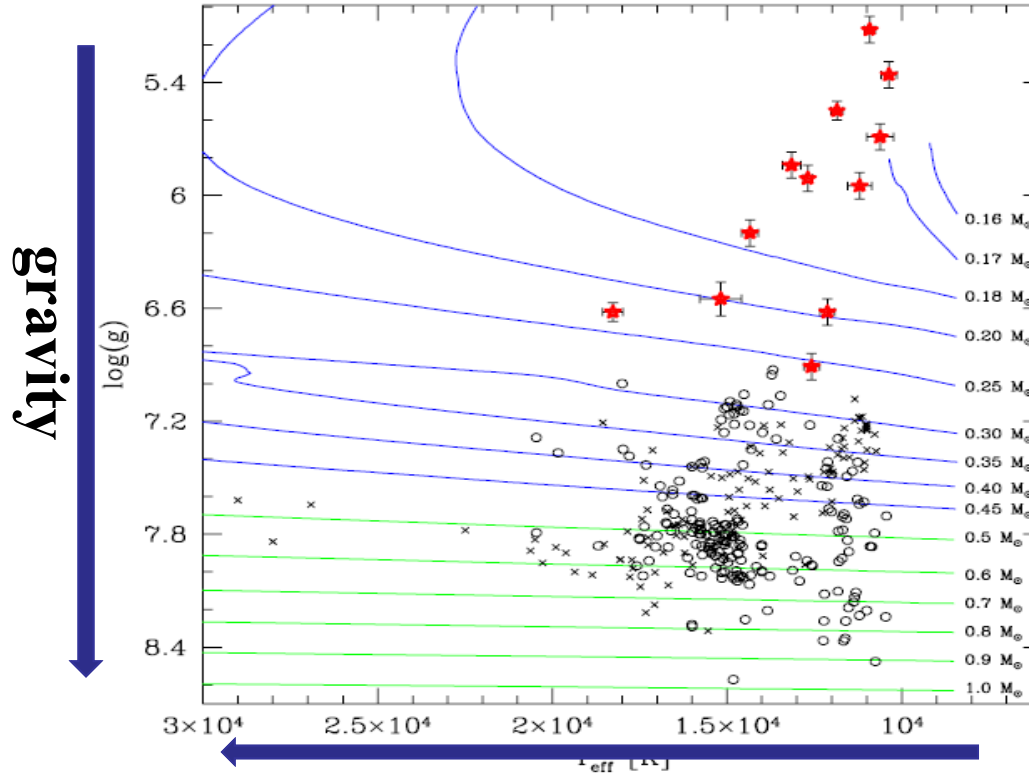


Outline

- Binary evolution
- **Binary-related objects**
- Evolutionary population synthesis
- Future perspectives

ELM WDs : Extremely low-mass white dwarfs

Brown+, 2010, ApJ



$5 < \log g < 7$

$8000 \text{ K} < T_{\text{eff}} < 22000 \text{ K}$

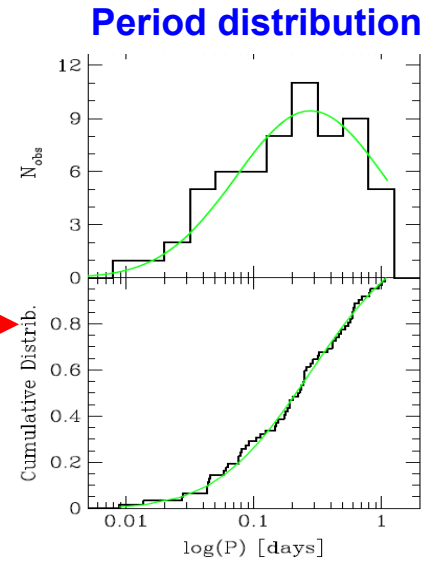
$M < 0.25\text{-}0.3 M_{\odot}$

Observations

I. The ELM Survey (Brown, Kilic et al. 2009--, SAO, 6.5 m MMT)

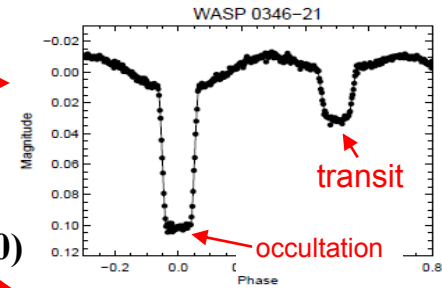
- 88 ELM WDs, of which 76 being in binaries
- Half of the observed binaries will merge in less than 6Gyr

Brown et al. 2016, ApJ



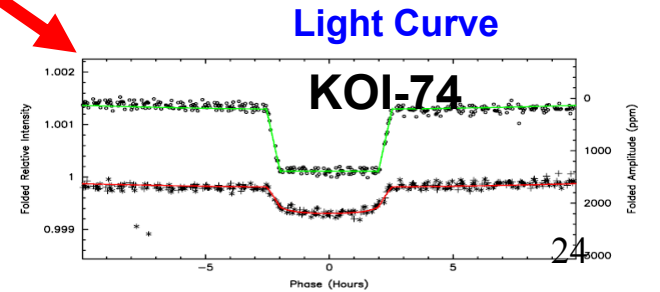
II. The WASP Project Maxted et al.

- The 1st sample: 1SWASP J024743.37-251549.2 (2011 MNRAS)
- Multiple-period pulsations (2013, Nature)
- 17 EL CVn-type binaries (2014, MNRAS)



III. Kepler

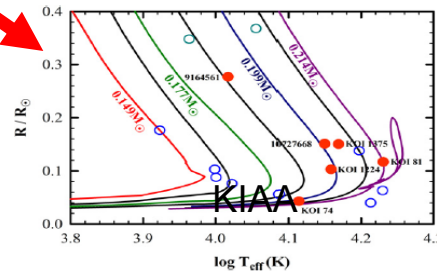
- The first two: KOI-74, KOI-81 (van Kerkwijk+,2010)
- 7 objects in total (Rappaport+,2015, Guo+, 2017)



VI. ELM in MSPs (>10)

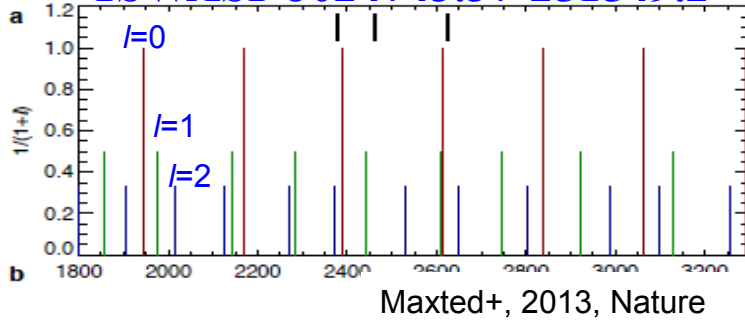
Istrate+, 2014, 2016

22-Mar-2018

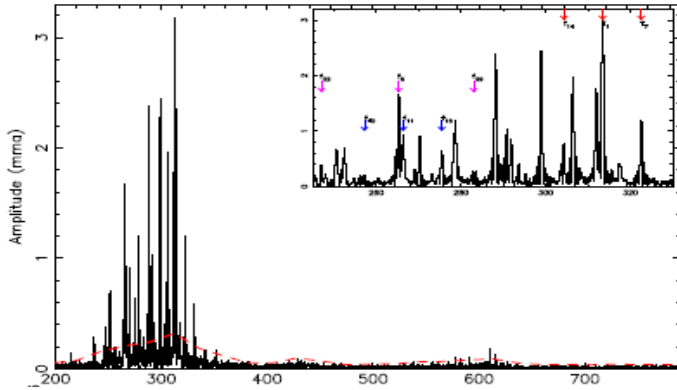


The pre-ELM WD instability strip

1SWASP J024743.37-251549.2

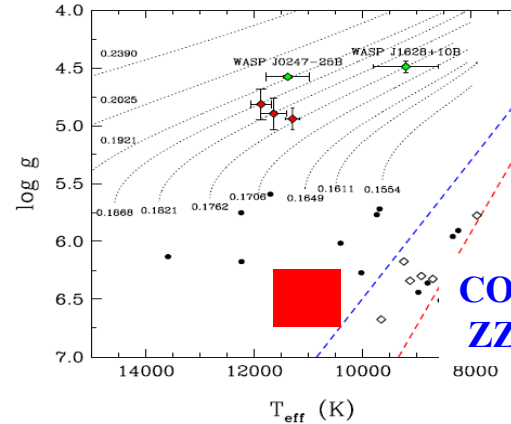


KIC 9164561



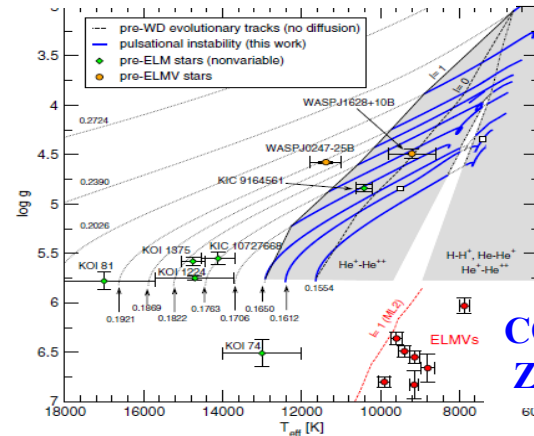
Zhang, X., 2016, ApJ

Diven by He⁺⁺, He⁺, H⁺



Corsico+, 2016, A&A

CO-core
ZZ Ceti stars

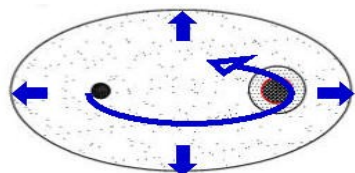


CO-core
ZZ Ceti stars

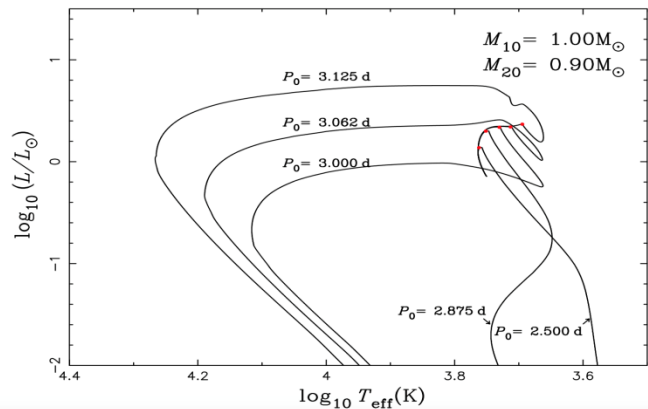
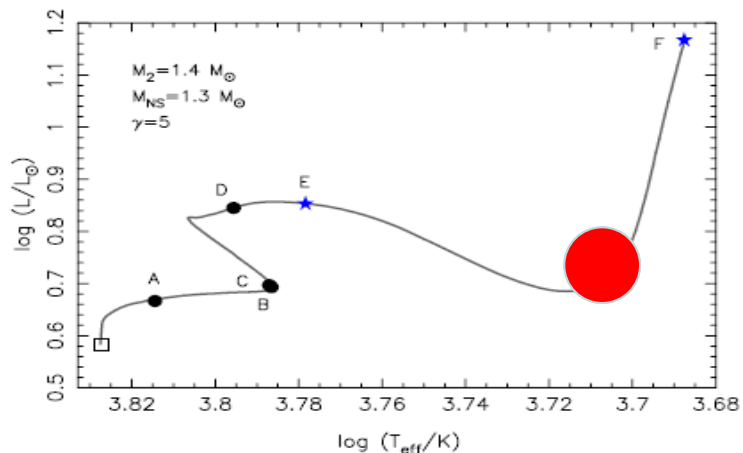
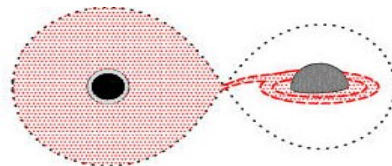
Formation of extremely low-mass WDs

Binary Evolution

common envelope ejection



stable mass transfer

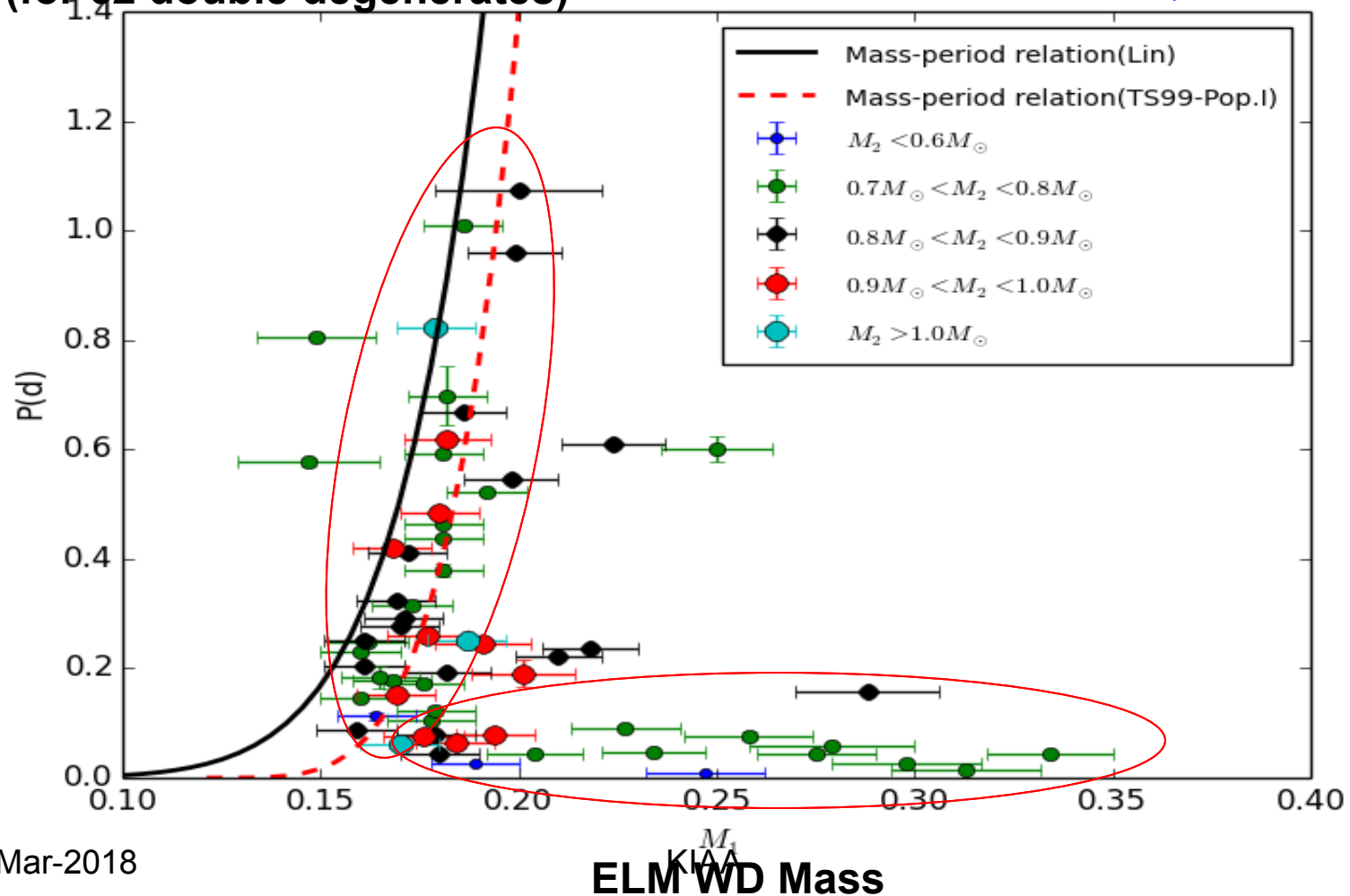


Beyond but very close to the bifurcation period

ELM WD mass VS orbital period

(for 62 double degenerates)

Z. Li,+

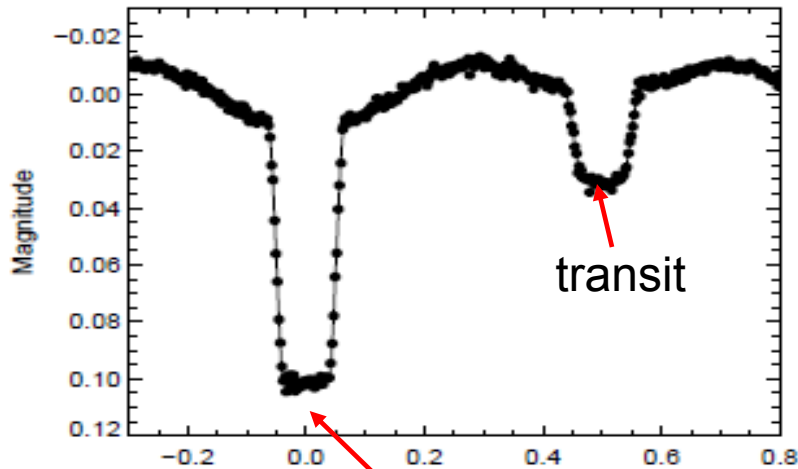


EL CVn-type Binaries

A/F dwarf star + proto-He WD

Maxted et al. 2014 , MNRAS

WASP 0346-21



The hot small companion
has been occulted.

The hot small companion:

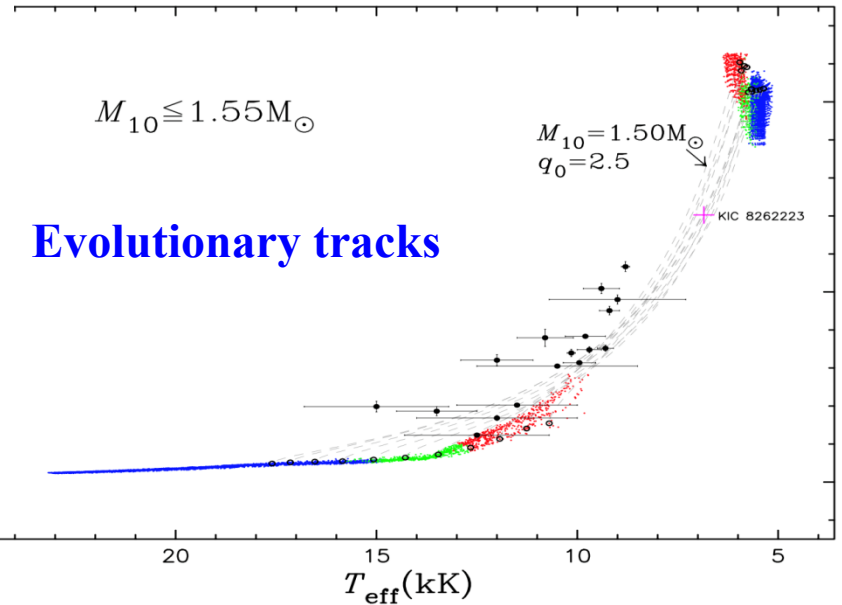
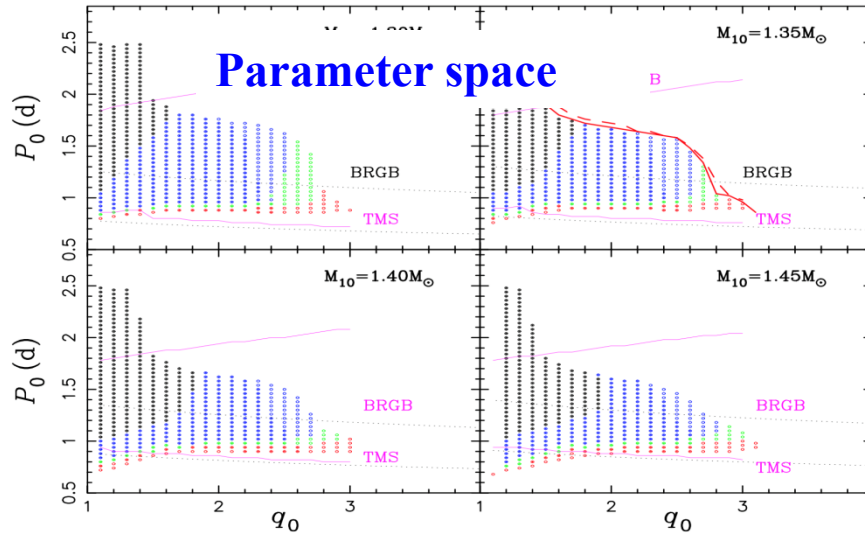
➤ **extremely low mass ($\sim 0.2M_{\text{sun}}$)**

➤ **in an evolutionary phase that
should be rare observed**

17 EL CVn-type binaries have been found
by Maxted et al. from the WASP observations.

The formation of EL CVn-type binaries

Chen X. et al 2017 , MNRAS, 464, 187



Predictions:

- More EL CVn-type binaries will be found in the Galaxy.
- They are more likely in the thin disk.

22-Mar-2018
(The peak mass for the progenitors is in $\sim 1.1-1.3 M_{\text{sun}}$)

KJAA

Discovery of 36 eclipsing EL CVn binaries found by the Palomar Transient Factory

J. van Roestel,¹★ T. Kupfer,² R. Ruiz-Carmona,¹ P. I. Groot,¹ T. A. Prince,² K. Burdøe,²
R. Laher,³ D. L. Shupe⁴ and E. Bellm⁵

¹Department of Astrophysics (MAAP), Durham University, Leazes Road, Durham, DL2 8TA, UK

Introduction for the model

In the canonical formation channel of EL CVn binaries (e.g. Chen et al. 2017), two main-sequence stars of similar mass are born at a short orbital period of a few days. The more massive star evolves faster and increases in radius. Before it ascends the red giant branch (RGB), it fills its Roche lobe and starts stable mass transfer to the lower-mass secondary star. This process continues until almost the complete outer envelope is transferred (identified as R CMa-type binaries; e.g. Lee et al. 2016). The remnant of the initially more massive star has become a pre-WD with a helium core and a thick hydrogen envelope ($\approx 0.01\text{--}0.04 M_{\odot}$; see Istrate et al. 2016a; Chen et al. 2017). The accretor has become a rejuvenated main-sequence

Transient Factory (PTF), combined with Sloan Digital Sky Survey (SDSS), Panoramic Survey Telescope and Rapid

of EL CVn systems. We do find tentative results that EL CVn systems occur more often in the thin disc, as was predicted by Chen et al. (2017). We also find that the space density is at the lower bound, or even below the prediction of $4\text{--}10 \times 10^{-6} \text{pc}^{-3}$. To prop-

Predictions of the model

Using stellar evolution and population synthesis codes, Chen et al. (2017) predict a space density of $4\text{--}10 \times 10^{-6} \text{pc}^{-3}$ for EL CVn binaries (including non-eclipsing ones) with orbital periods less than 2.2 d. In addition, they predicted that EL CVn binaries should mainly be found in young stellar populations, and therefore should be more abundant in the thin disc. We use the Galaxy model based

Space density : in line with Chen et al.

systems contain small pre-WDs ($<0.05 M_{\odot}$) and, as a consequence, their light curves feature shallow eclipses only detectable from space. The fact that 10 EL CVn-like systems are found in the Kepler field suggests that there should be many more in our Galaxy, in line with an estimate of the local space density from stellar evolution and population synthesis models, $4\text{--}10 \times 10^{-6} \text{pc}^{-3}$ (Chen et al. 2017).

ambiguous cases are from the thin disc, this probability drops well below 1 per cent. This indicates that our model is unlikely to be correct, and confirms that EL CVn systems are more abundant in the thin disc compared to the average stellar population, as was already suggested by Chen et al. (2017).

KIC 8262223 Z. Guo, D. Gies, R. Mstson et al (2017, ApJ)

KIC 8262223: A Post-mass Transfer Eclipsing Binary Consisting of a Delta Scuti Pulsator and a Helium White Dwarf Precursor

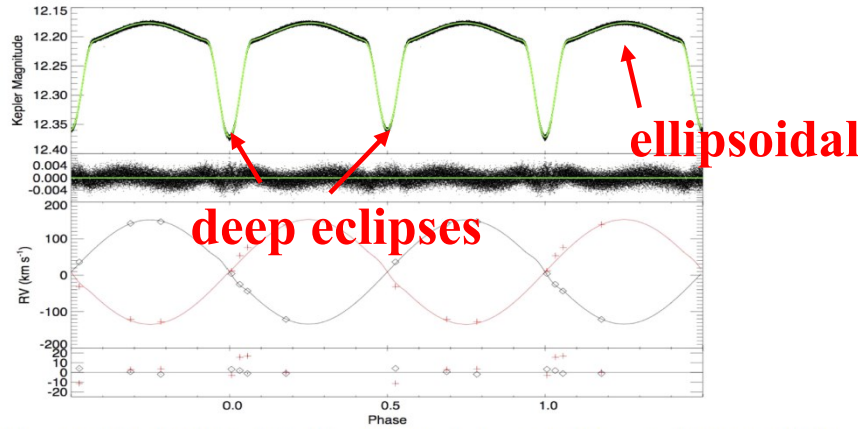


Figure 2. The upper and lower panels show the primary and secondary star a

delta Scuti pulsation

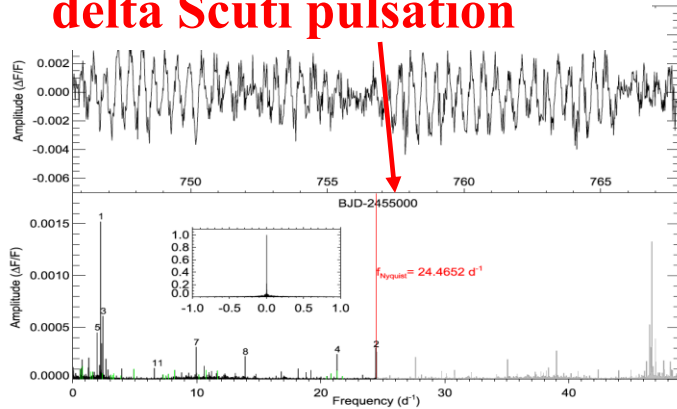
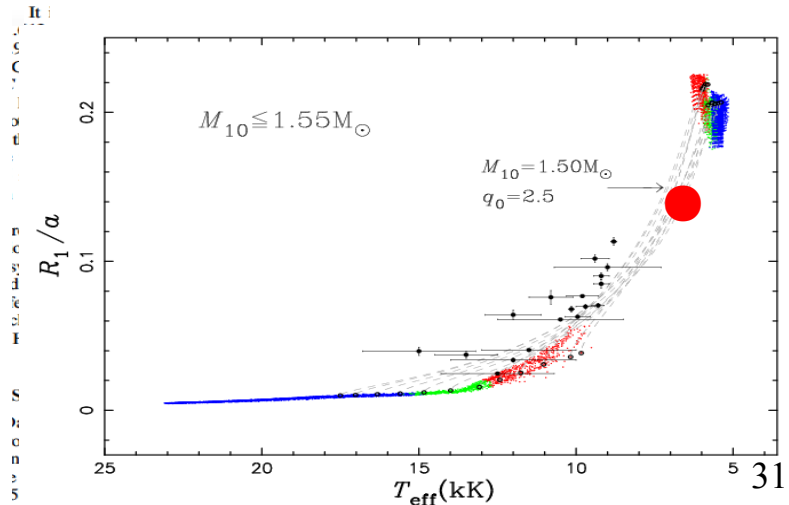
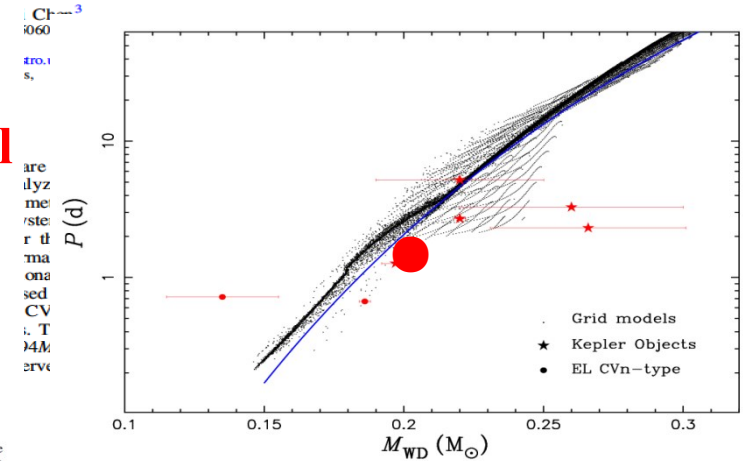


Figure 3. Time series of *Kepler* photometric data of oscillations with the binary light curve subtracted in the time domain (upper) and frequency domain (lower). Several of the strongest pulsation peaks have been labeled by their frequency numbers (Tables 3 and 4). The inset in the lower panel shows the spectral window. The Fourier spectrum in the super-Nyquist region is also shown, with the red vertical line indicating the Nyquist frequency $f_{\text{Nyquist}} = 24.4652 \text{ day}^{-1}$ of *Kepler* long-cadence data. All pulsations in the super-Nyquist region (gray shaded) have lower amplitudes compared to their mirror reflections around the f_{Nyquist} , indicating they are not real pulsation peaks, but are aliases.

secondary ($m = 0.21 M_{\odot}$, $R = 1.13 R_{\odot}$) and a δ Scuti pulsating primary.

KIC 8262223, $K = 12.146$, $\text{mag} = 20:01:19.788$, $\text{mag} =$

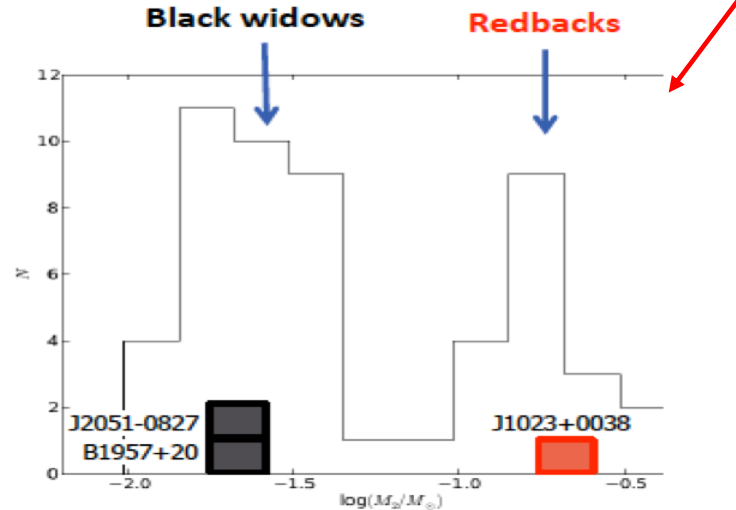
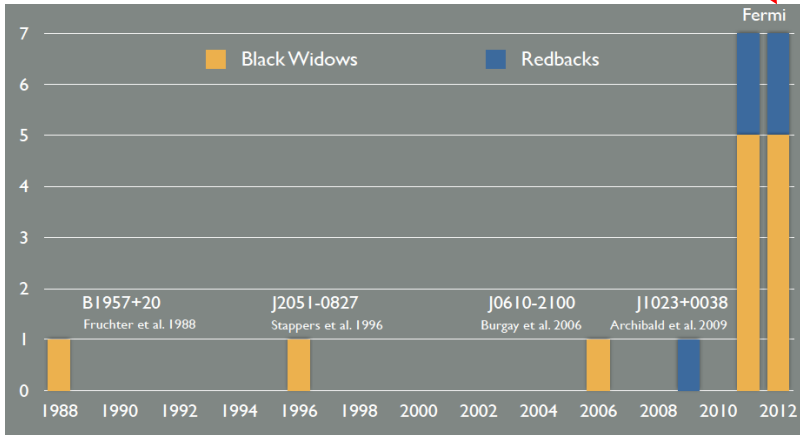


minutes, respectively. Please refer to Caldwell et al. (2017) for more information on the *Kepler* data. The aperture contamination factors (k) reported in the *Kepler* Input Catalog are lower

Eclipsing Binary Millisecond Pulsars

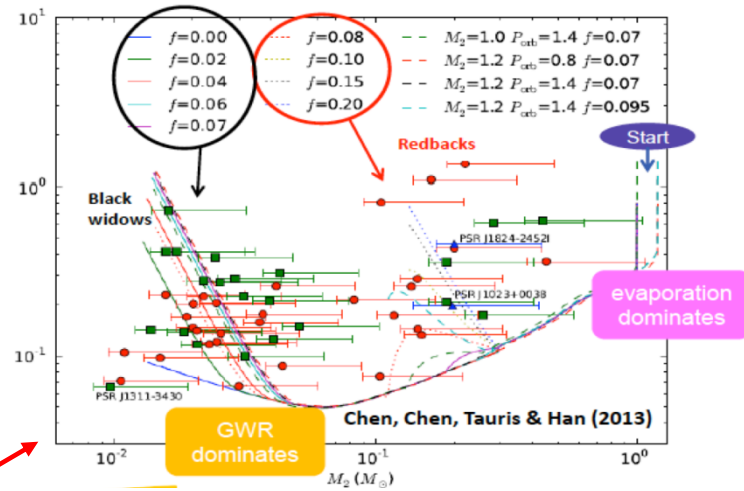
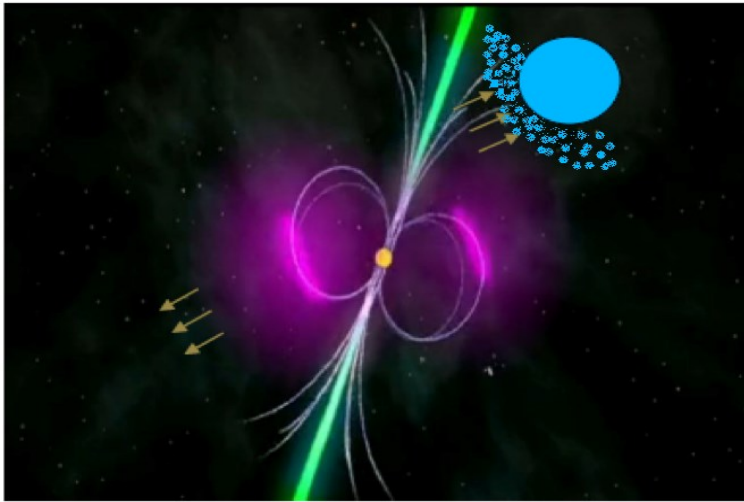
Discovered by Fermi

Two distinct populations



- The masses of the companions are very low, then they should be small. **How to produce the eclipses ?**
- Both populations have orbital periods between 0.1 and 1.0 days, **their companion masses differ by an order of magnitude. Why?**

➤ **Pulsar irradiation and evaporation of the companions**
are crucial for the formation of eclipsing MSPs



➤ The determining factor for the producing the two populations is **the efficiency of the irradiation process.**

widely accepted in this field

Dick Manchester

《Physics of neutron stars and Related objects》 (2017, JApA, 38, 42) :



Millisecond Pulsars, their Evolution and Applications

R. N. Manchester

CSIRO Astronomy and Space Science, PO Box 76, Epp
dick.manchester@csiro.au

28 September 2017

ABSTRACT

Millisecond pulsars
mal” pulsars, not c
rates and high prob
with MSPs having
star formation in an

Another population of binary systems with very short orbital periods ($P_b < 1$ day) but somewhat larger companion masses, mostly between 0.1 and 0.3 M_\odot , can be identified on the $m_c - P_b$ plot. These are the other type of “spider” pulsar known as redbacks (Roberts 2013). The companion is a non-degenerate star that fills its Roche lobe and is losing mass, partly as a result of heating by the pulsar wind, resulting in extended and variable eclipses of the pulsar radio signal – see Chen et al. (2013) for a discussion of the formation processes of black-widow and redback eclipsing binaries. The

any globular cluster and, hence, an origin in a cluster is very unlikely. This showed that such systems could result from evolution of a binary system without the need to invoke exchange interactions; possible evolutionary paths are discussed by Chen et al. (2013). Since then, the *Fermi*-related

the MSPs nearly ideal probes of a wide , they have been used to detect plan- tational theories, to set limits on the the Universe, and to establish pulsar- ck timescales in long-term stability. ary evolution, often suggesting exotic retion-powered MSPs, and especially MSP and a non-accreting radio MSP, give important insight into the physics of accretion on to highly magnetised neutron stars.

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Key words: pulsars: general—stars: evolution—gravitation

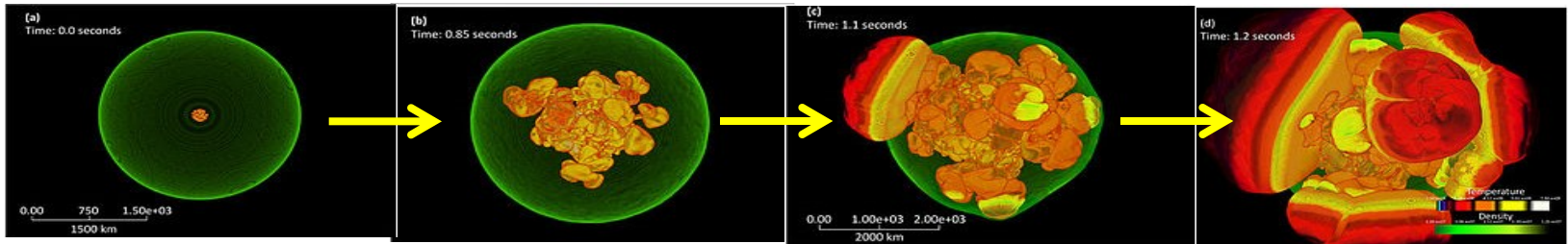
SNe Ia

Thermonuclear Explosions of CO WDs at $\sim 1.4 M_{\text{sun}}$

carbon ignition under degenerate conditions \longrightarrow
thermonuclear runaway \longrightarrow

incineration and complete destruction of the star

10^{51} ergs



Phillips Relation

(Phillips, 1993): empirical

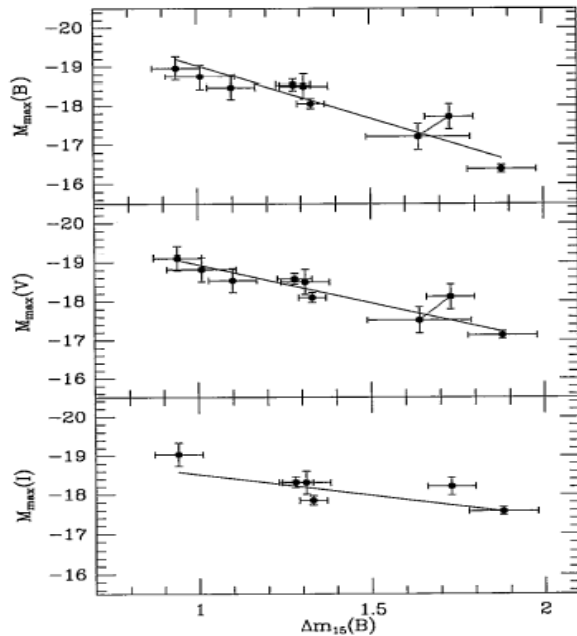
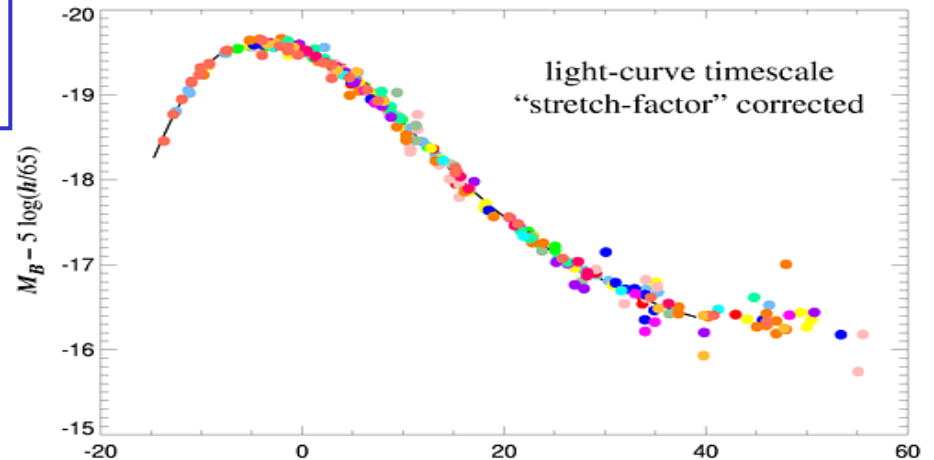
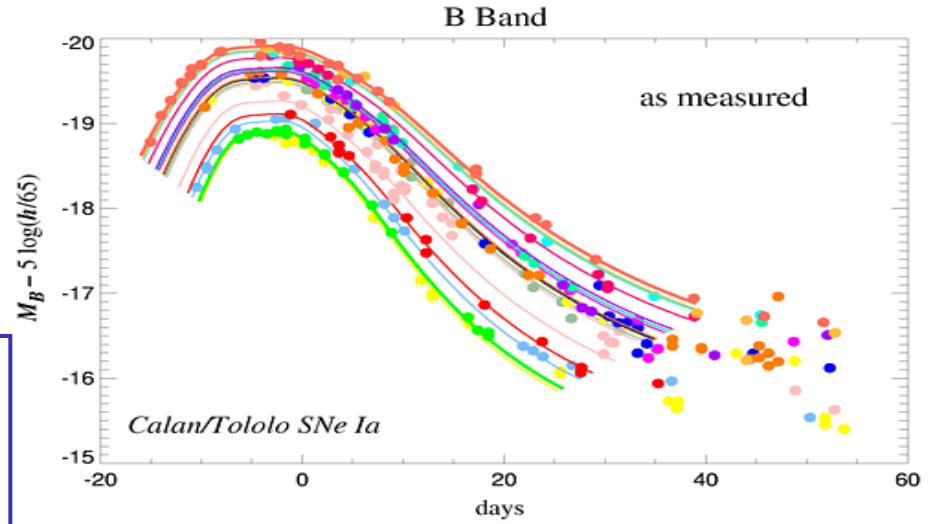


FIG. 1.—Decline rate–peak luminosity relation for the nine best-observed SN Ia's. Absolute magnitudes in B , V , and I are plotted vs. $\Delta m_{15}(B)$, which measures the amount in magnitudes that the B light curve drops during the first 15 days following maximum.

22-Mar-2018

Luminosity

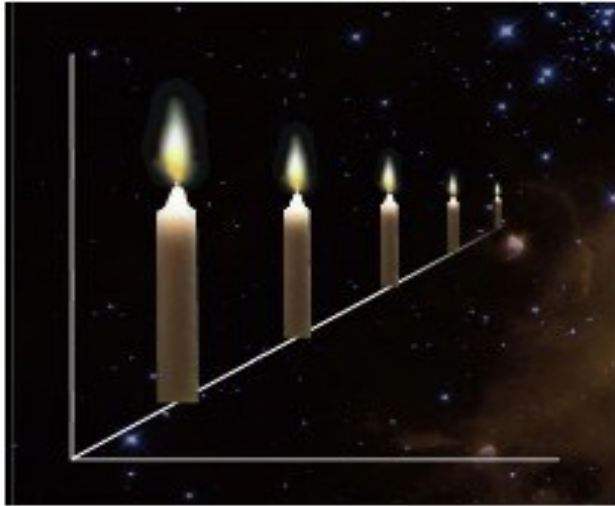


time

KIAA

Kim, et al. (1996)

SNe Ia as distance indicators



Accretion of WDs

The mass of CO WDs at birth is less than $\sim 1.1 M_{\text{sun}}$.
($\sim 0.6 M_{\text{sun}}$)



Accretion is necessary for final explosions .

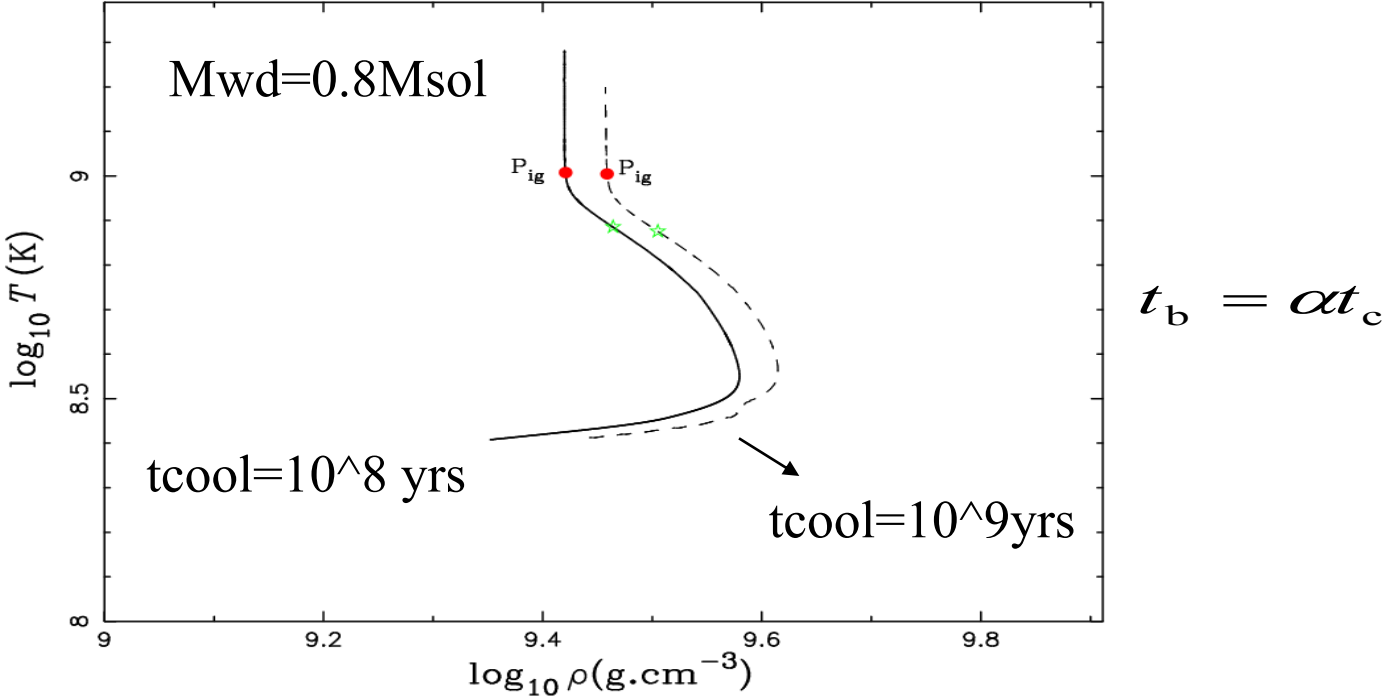
How to accrete ?

and when to explode ?

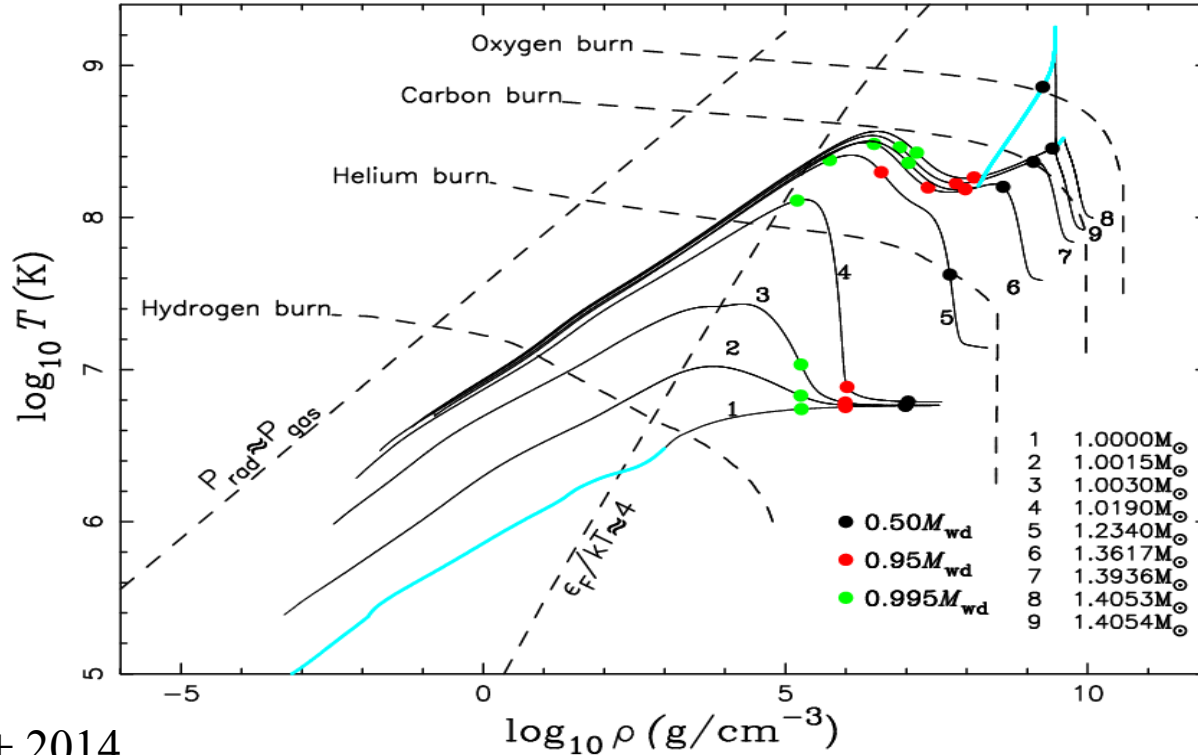
Maximum CO WD mass at birth: 1.1Msol

Chen+ 2014, MNRAS

When does C ignition occur ? $T_{\max} \uparrow$ while keeping ρ constant.



CO WD structure during accretion



Chen+ 2014

Progenitor models of SNe Ia

Single-degenerate (SD) channel:

accreting from non-degenerate companion

CO WD + MS

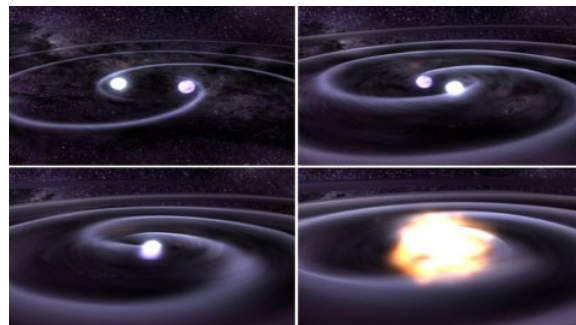
CO WD + RG

CO WD + He Star

Double-degenerate (DD) channel:

Two CO WDs with a combined mass
greater than $1.4M_{\text{sun}}$ merge

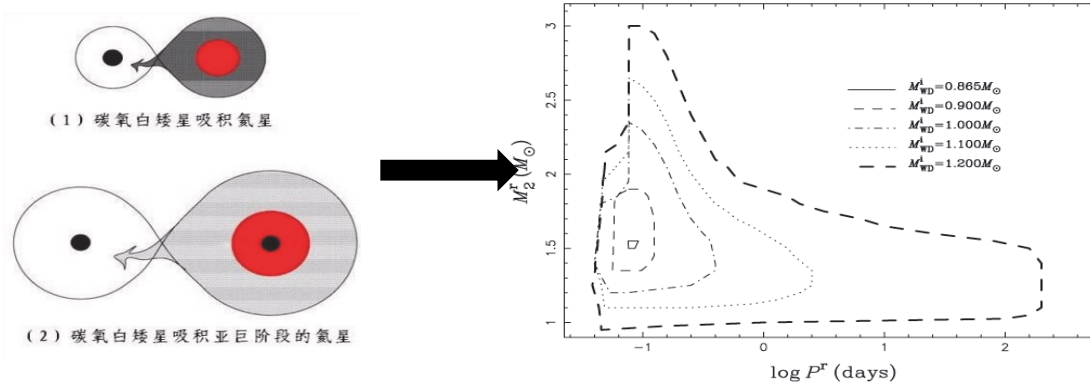
CO WD + CO WD



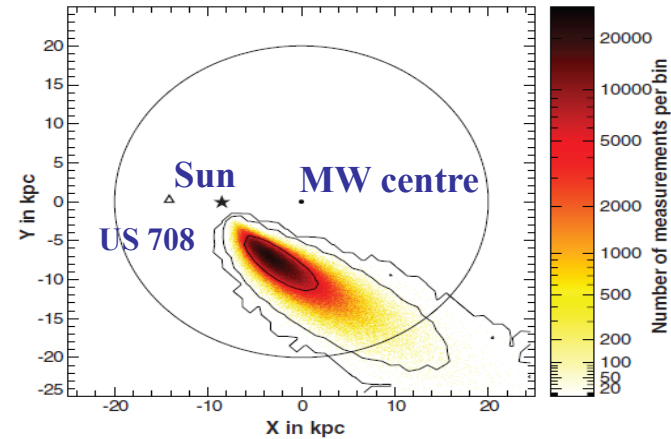
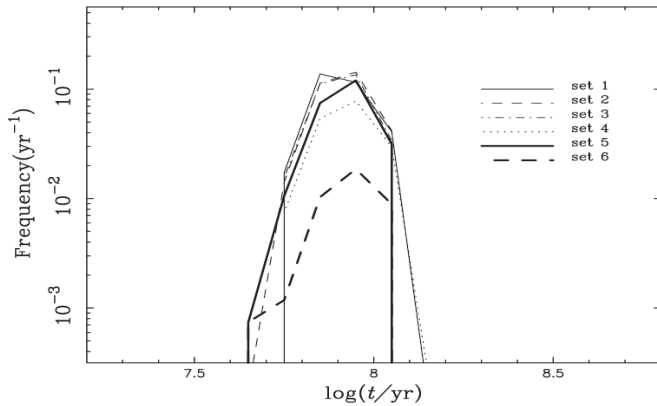
Credit: GSFC/Dana Berry

Helium donor channel (WD+He)

Wang+ 2009, 2010



Birthrate
Time delay
Properties of Progenitors
Properties of remnant stars
.....



Explain young SNe Ia

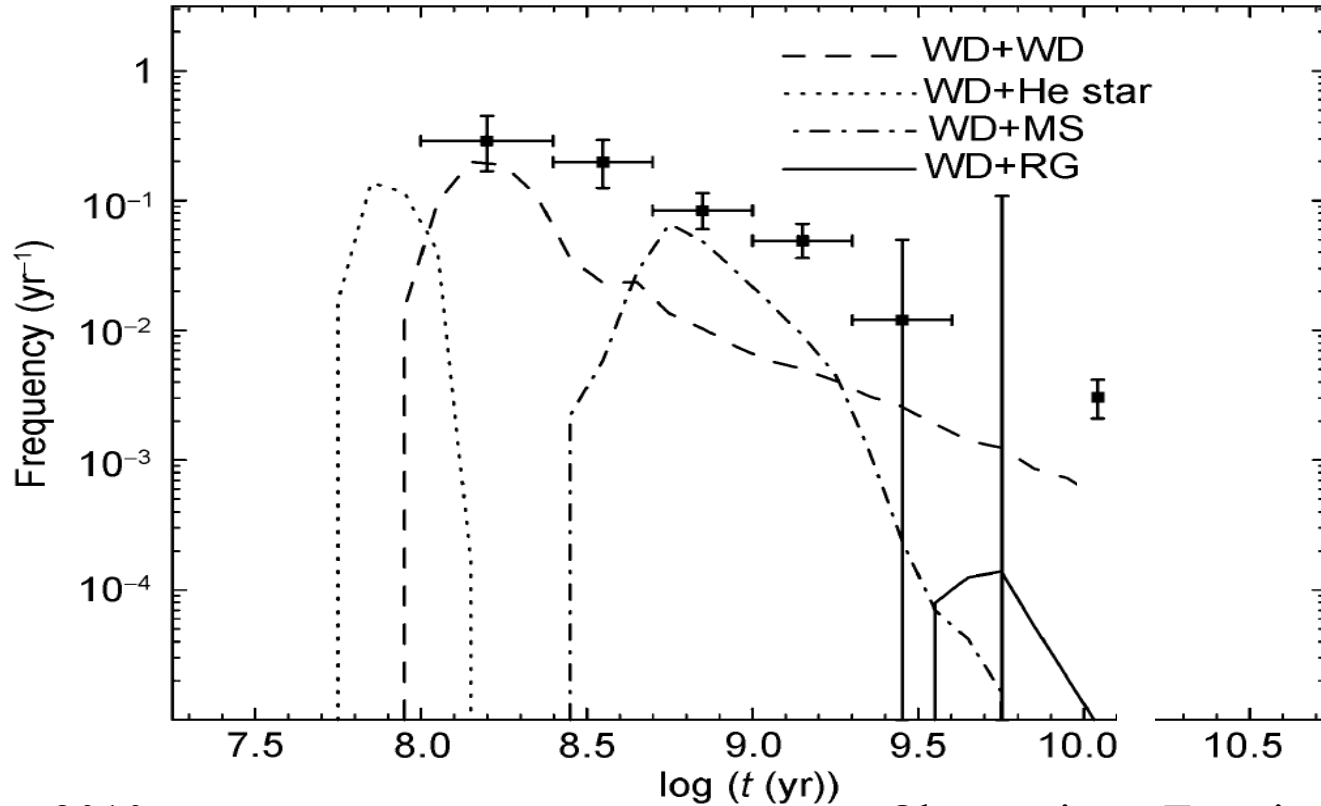
22-Mar-2018



The fastest unbound star in our Galaxy ejected by a thermonuclear supernova
S. Geier *et al.*
Science **347**, 1126 (2015); DOI: 10.1126/science.1259063

Hypervelocity star
Geier+, 2015, Science

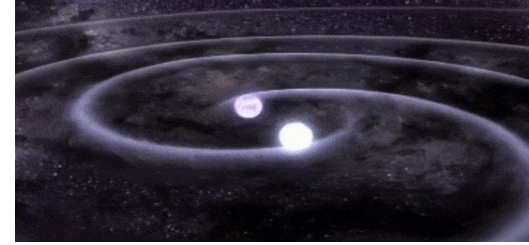
SN Ia rates for single star burst



Wang+ 2010

Observations: Totani+ 2008

Stellar mass double BHs



Massive binary evolution is the key scenario for their formation

CE evolution

$\sim 218 \text{ Gpc}^{-3}\text{yr}^{-1}$

Contact binary evolution

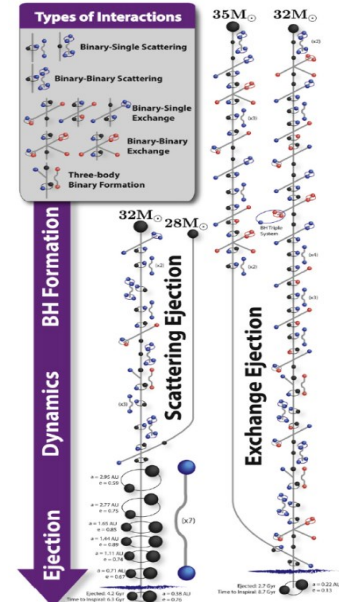
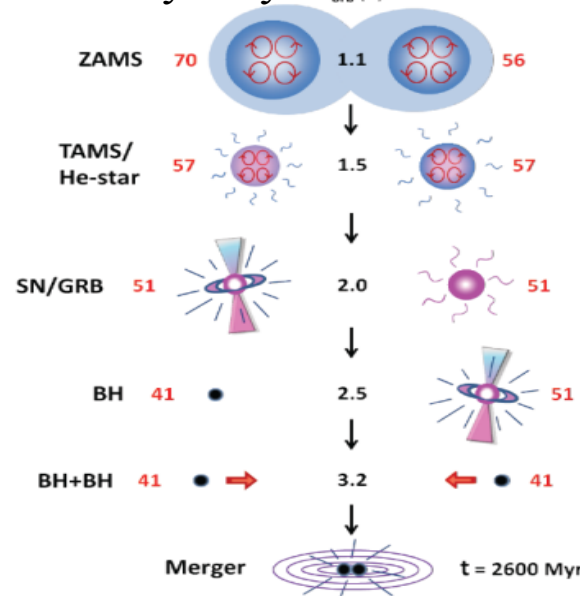
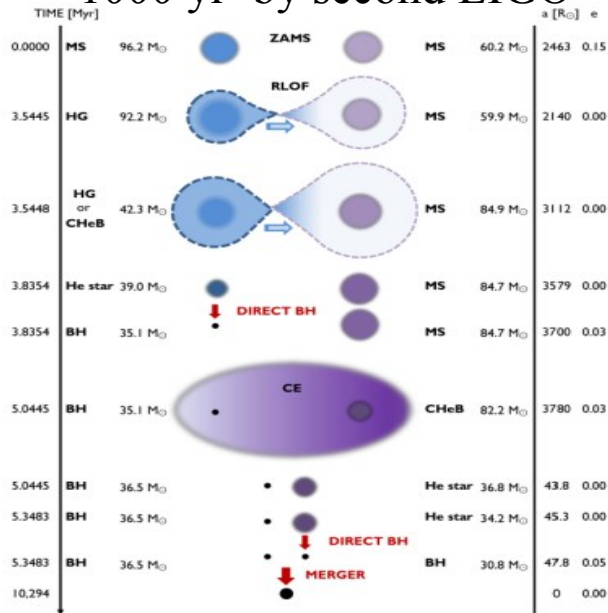
$\sim 10 \text{ Gpc}^{-3}\text{yr}^{-1}$

GC dynamics

$\sim 6.5 \text{ Gpc}^{-3}\text{yr}^{-1}$

$\sim 1000 \text{ yr}^{-1}$ by second LIGO

$\sim 500 \text{ yr}^{-1}$ by second LIGO



VFTS 352: 29 Msun+29 Msun with P=1.12 d

The formation of double degenerates and related objects

Zhanwen Han^{1,2}

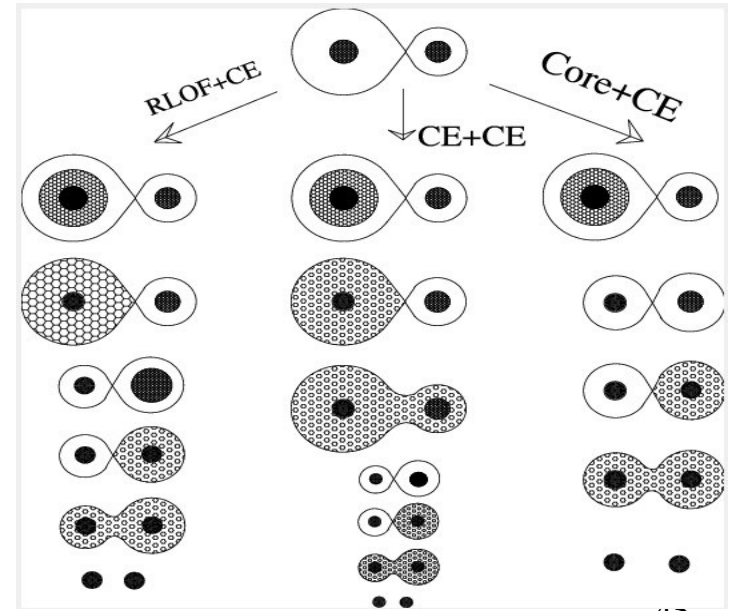
¹Centre for Astrophysics, University of Science and Technology of China, Hefei, 230026, China

²Yunnan Observatory, Academia Sinica, Kunming, 650011, China

Observation

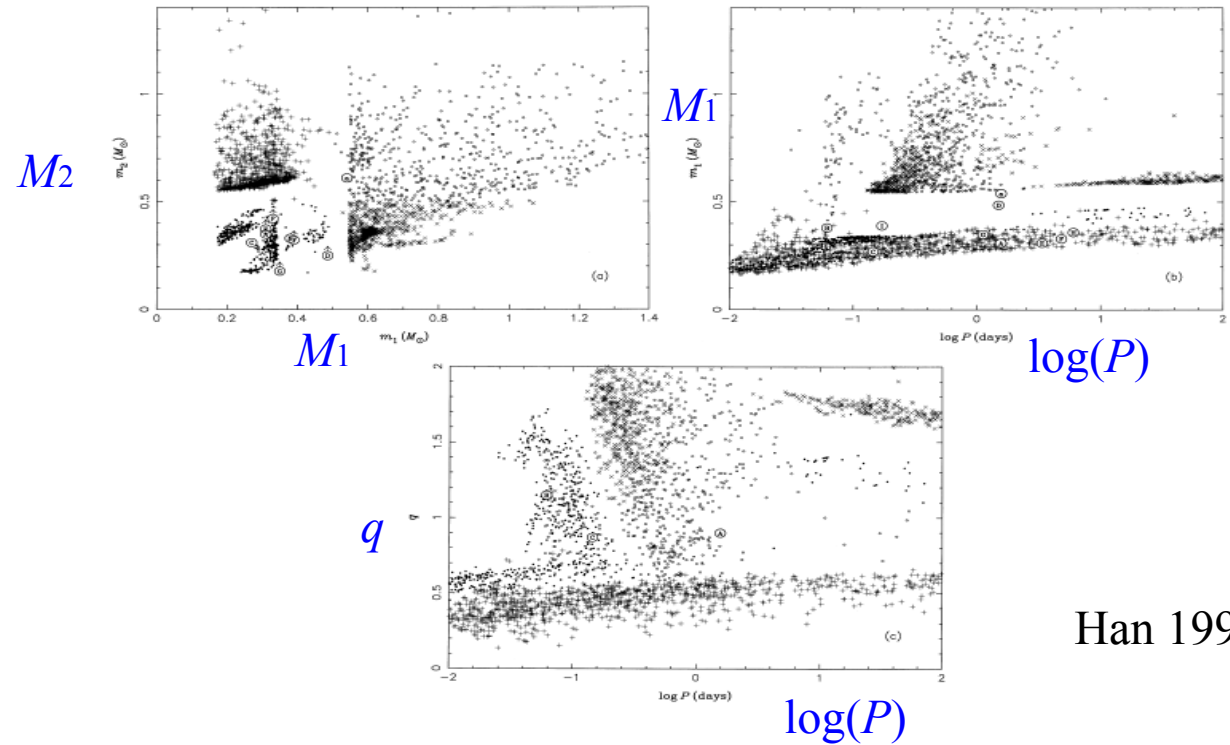
DD	WD number	Name	$M_1 (M_{\odot})$	$M_2 (M_{\odot})$ ([†] for lower limit)	Orbital Period (days)
A	0135–052	L870-2/EG11	0.31	0.35	1.56
a			0.54	0.61	
B	0957–666		0.38	0.33	0.061
C	1101+364	PG 1101+364	0.27	0.31	0.1446
D	1202+608	GD 314/LB 2197	0.486	0.25 [†]	1.49
E	1241–010	PG 1241-010	0.31	0.373 [†]	3.35
F	1317+453	G 177-31	0.33	0.421 [†]	4.8
G	1713+332	GD 360	0.35	0.178 [†]	1.12
H	2032+188	GD 231	0.36		2-10 (6 [†])
I	2331+290	GD 251	0.39	0.322 [†]	0.14-0.20 (0.17 [†])

Theory



Distributions of masses, mass ratios, orbital periods of Galactic DWDs with BPS approach

1.762×10^8 double WDs



Han 1998

Gravitational Radiation from Close Double White Dwarfs

Ronald F. Webbink* and Zhanwen Han†

* *Department of Astronomy, University of Illinois, Urbana, Illinois 61801, U.S.A.*

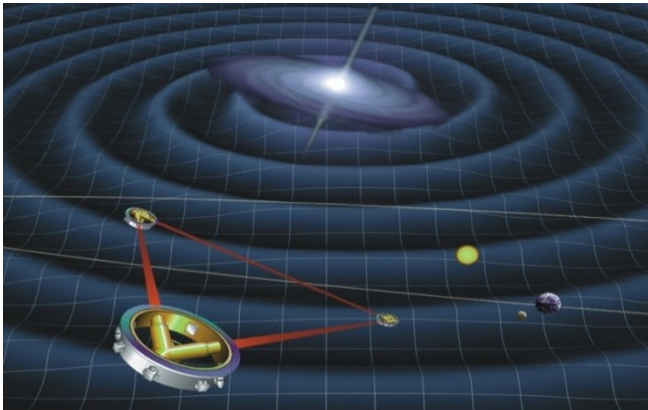
† *Centre for Astrophysics, University of Science and Technology of China, Hefei, 230026, China
and Yunnan Observatory, Academia Sinica, Kunming, 650011, China*

CP456, *Laser Interferometer Space Antenna*

edited by William M. Folkner

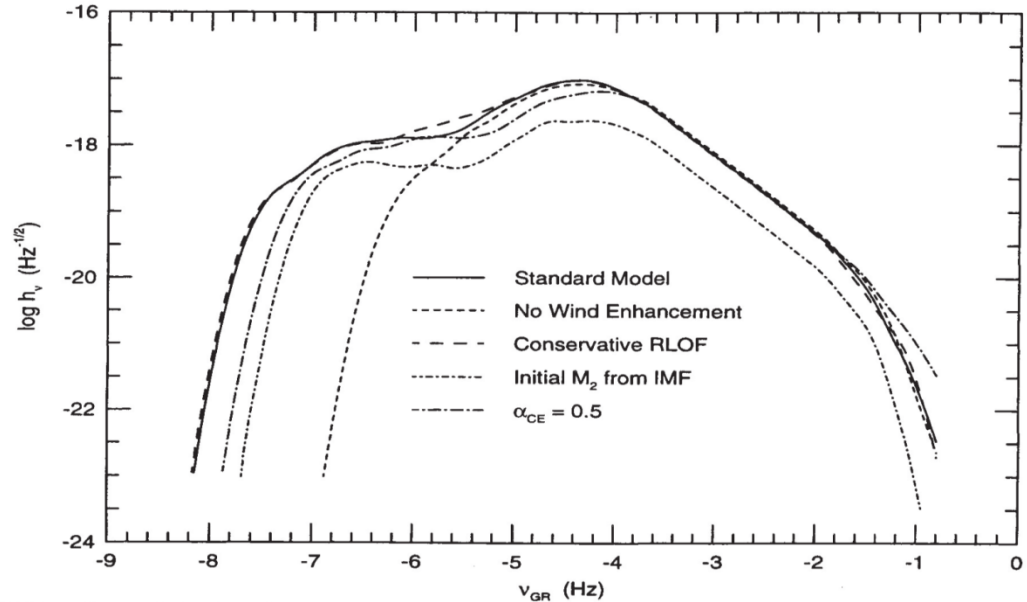
© 1998 The American Institute of Physics 1-56396-848-7/98/\$15.00

61

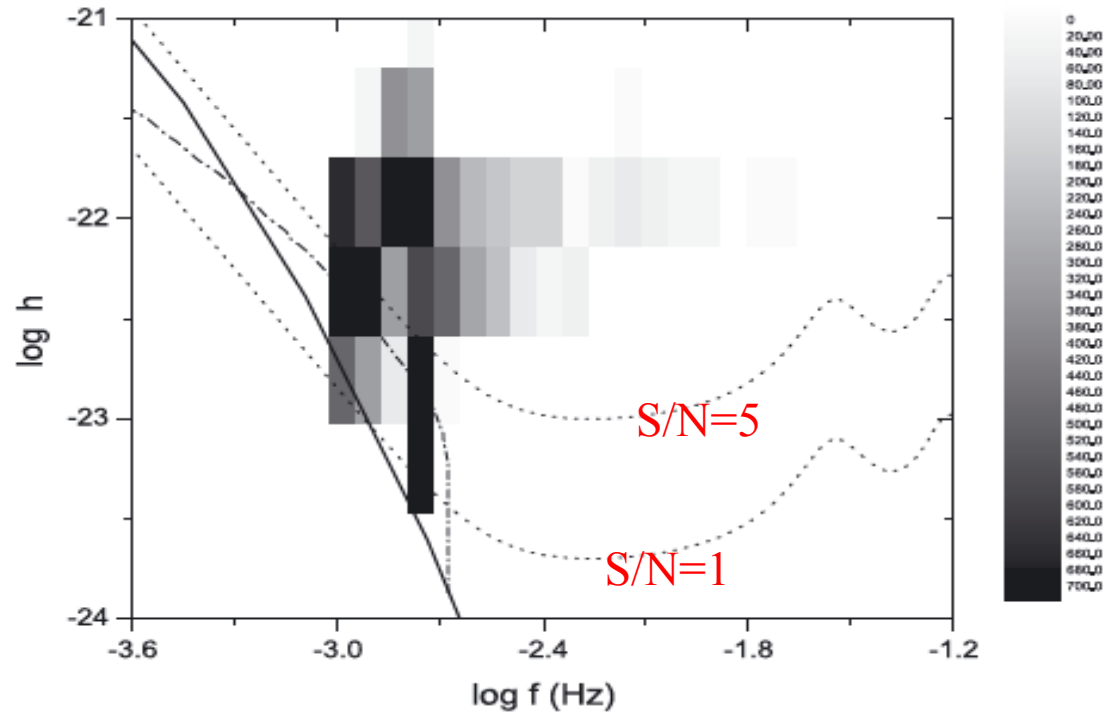


LISA

Spectral density of GW amplitude

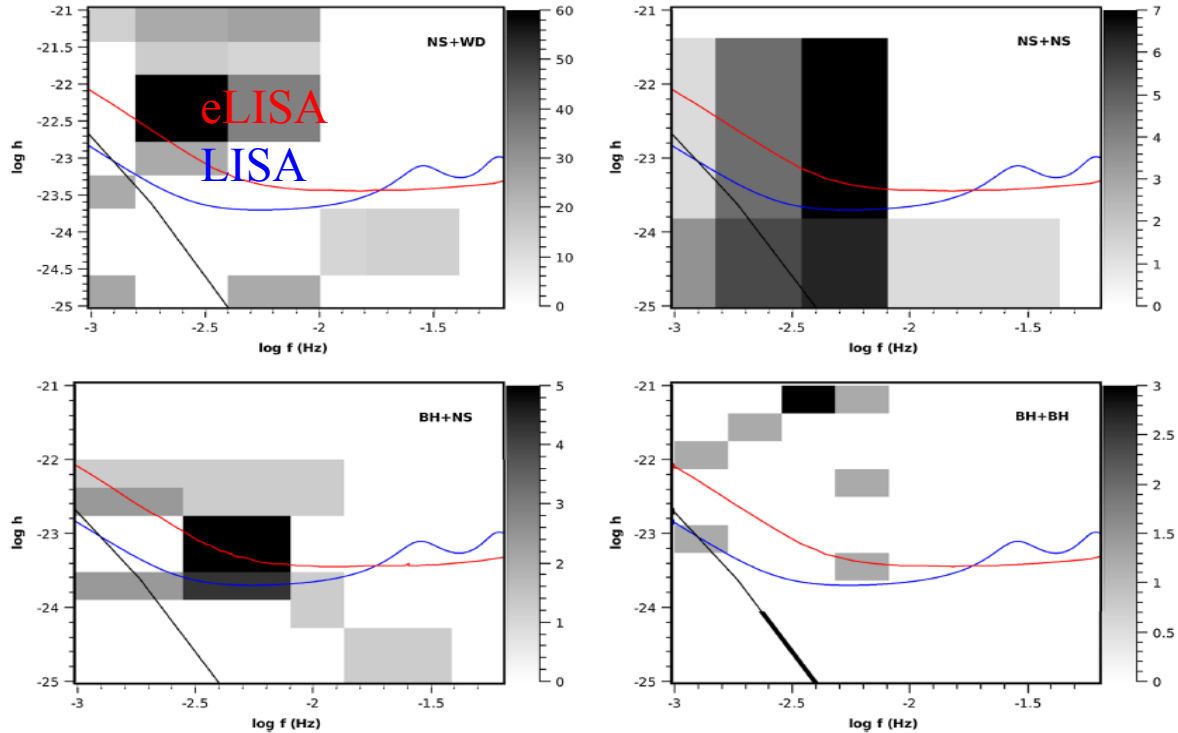


Resolvable DWDs with LISA

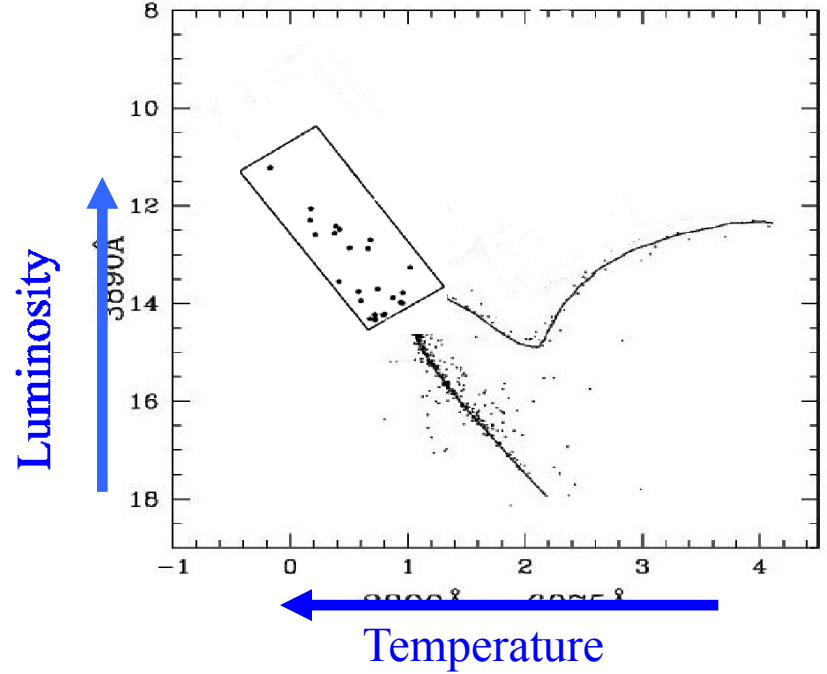
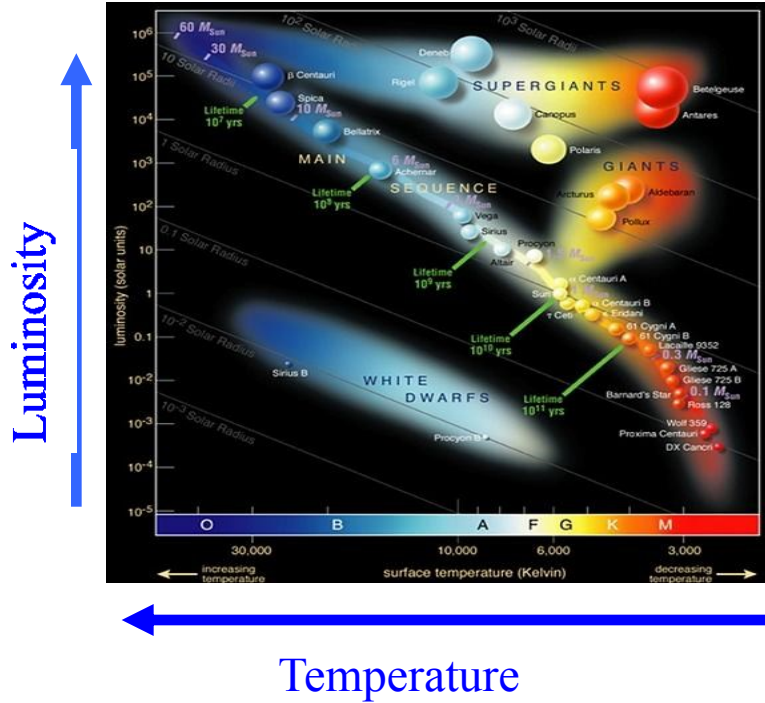


Liu+ 2009, ApJ

Resolvable NS+WD, NS+NS, BH+NS, BH+BH with eLISA

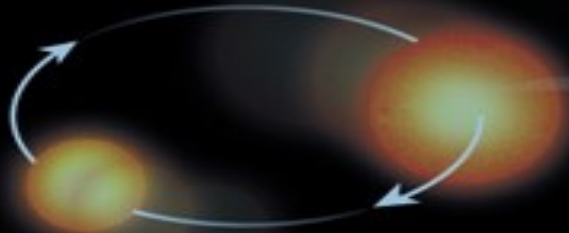


Blue stragglers



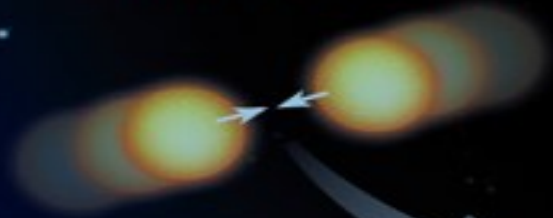
HR diagram

Stellar cannibalism



Chen+

Stellar collision



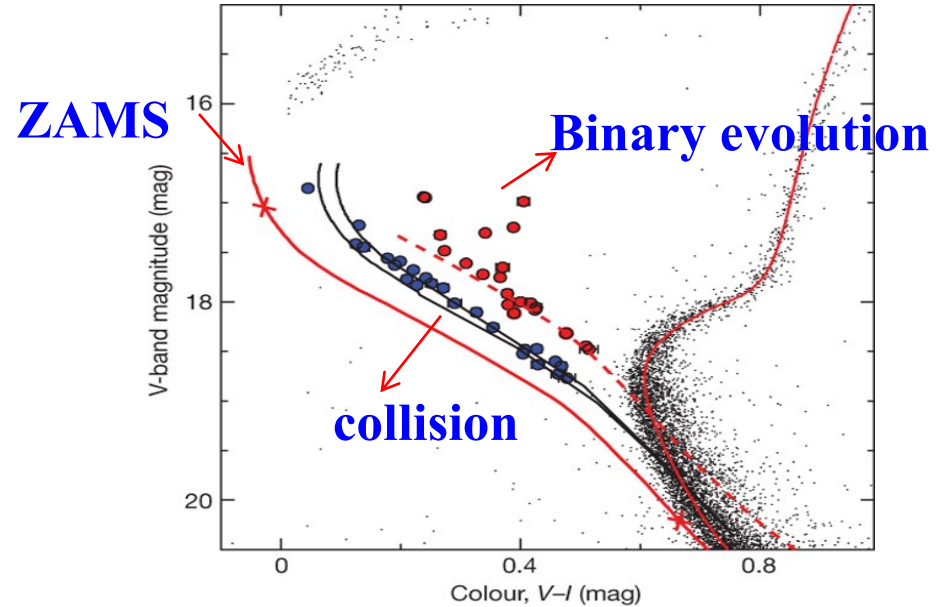
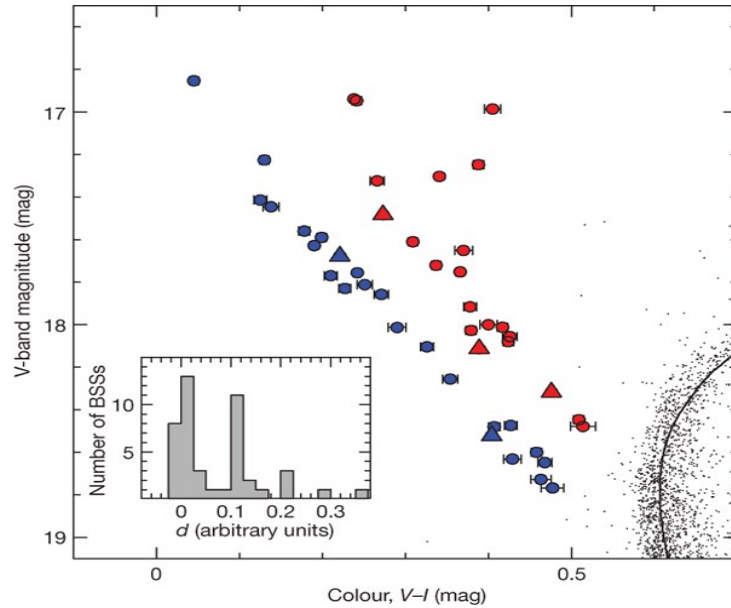
Blue straggler star

Sills+

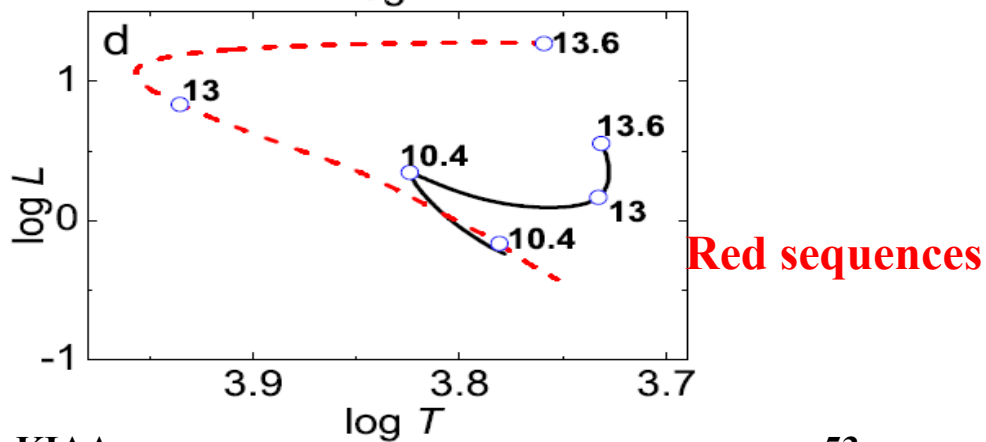
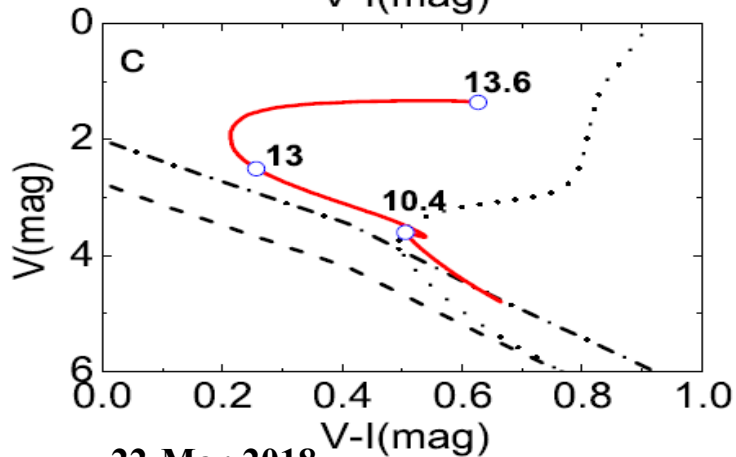
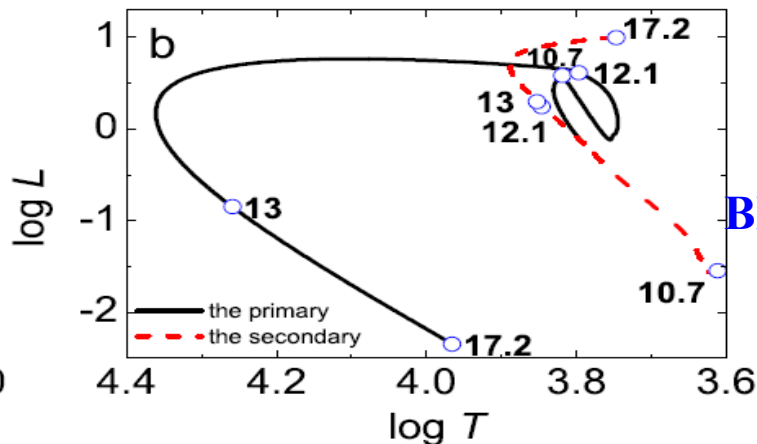
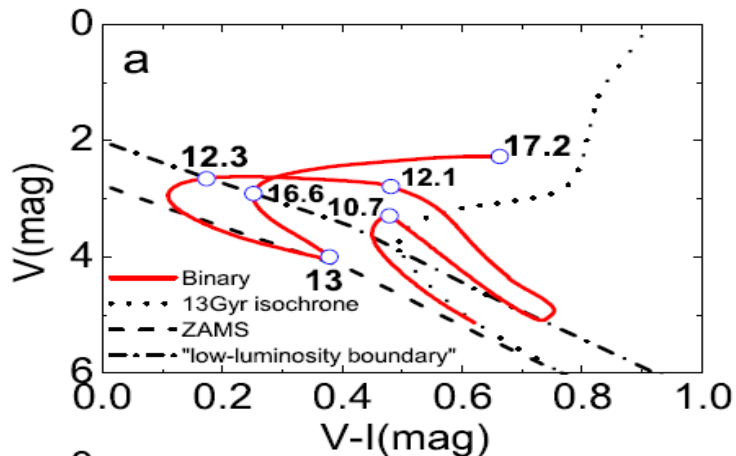


Two blue stragglers sequences in M30

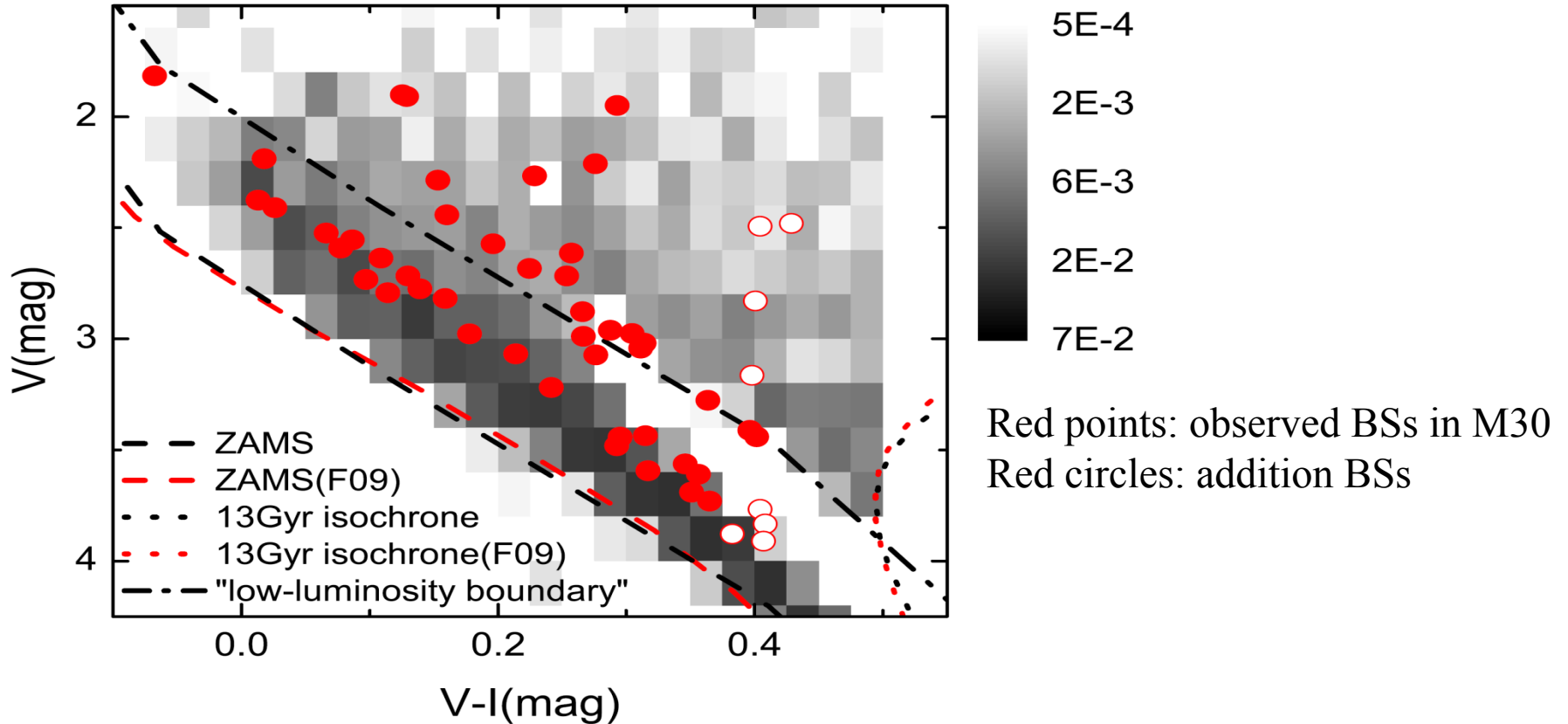
Ferraro+, 2009, Nature



Binary evolution may produce both of the sequences

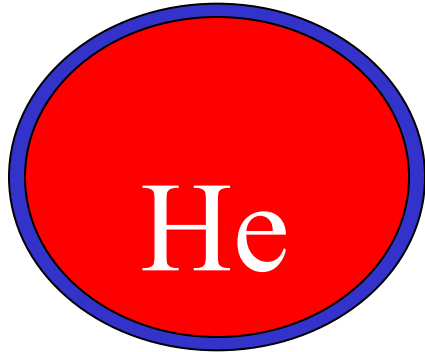


Blue stragglers with BPS



Jiang+, 2017, ApJ

hot subdwarf stars



$$M \sim 0.5 M_{\odot}$$

$$M_{\text{env}} \leq 0.02 M_{\odot}$$

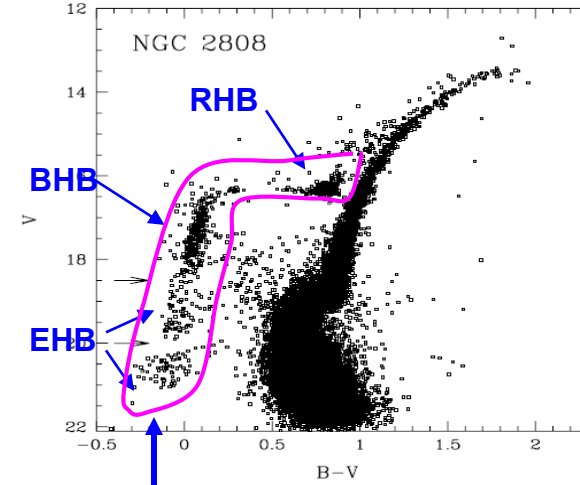
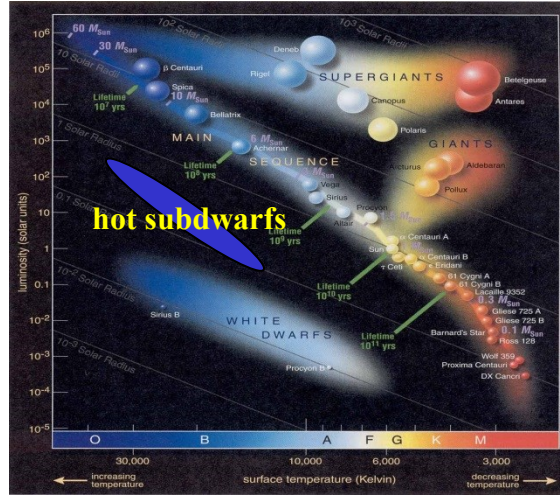
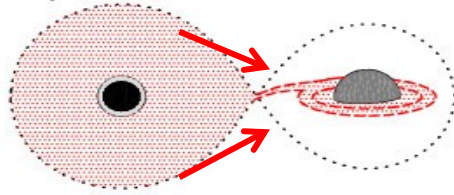


Fig. 4. The V vs. $(B - V)$ CMD from the *HST* WFPC2 data of the inner core of NGC 2808 ($r \leq 100$ arcsec). Only the 35,000 stars with a photometric error less than 0.15 are plotted.

Extreme Horizontal Branch stars

Binary model: stable RLOF

Stable RLOF near the tip of FGB



Wide hot subdwarf binary with MS companion

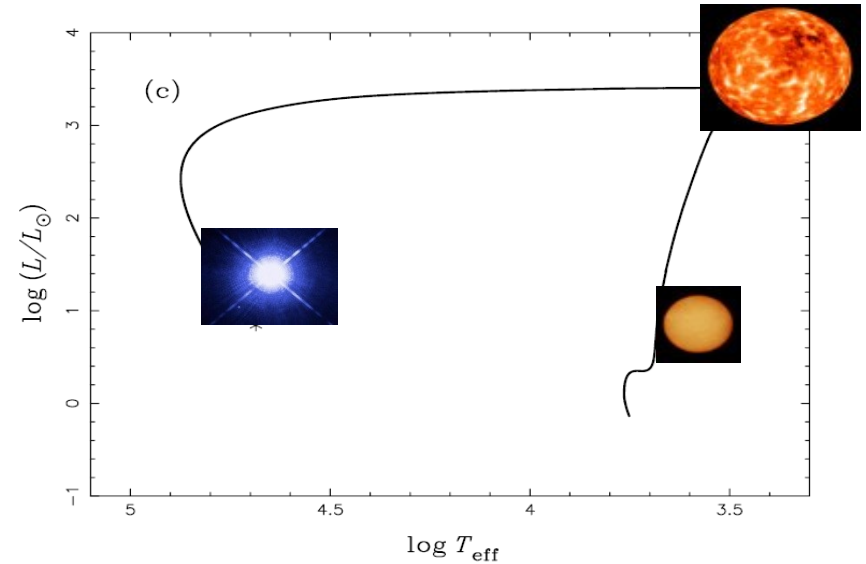
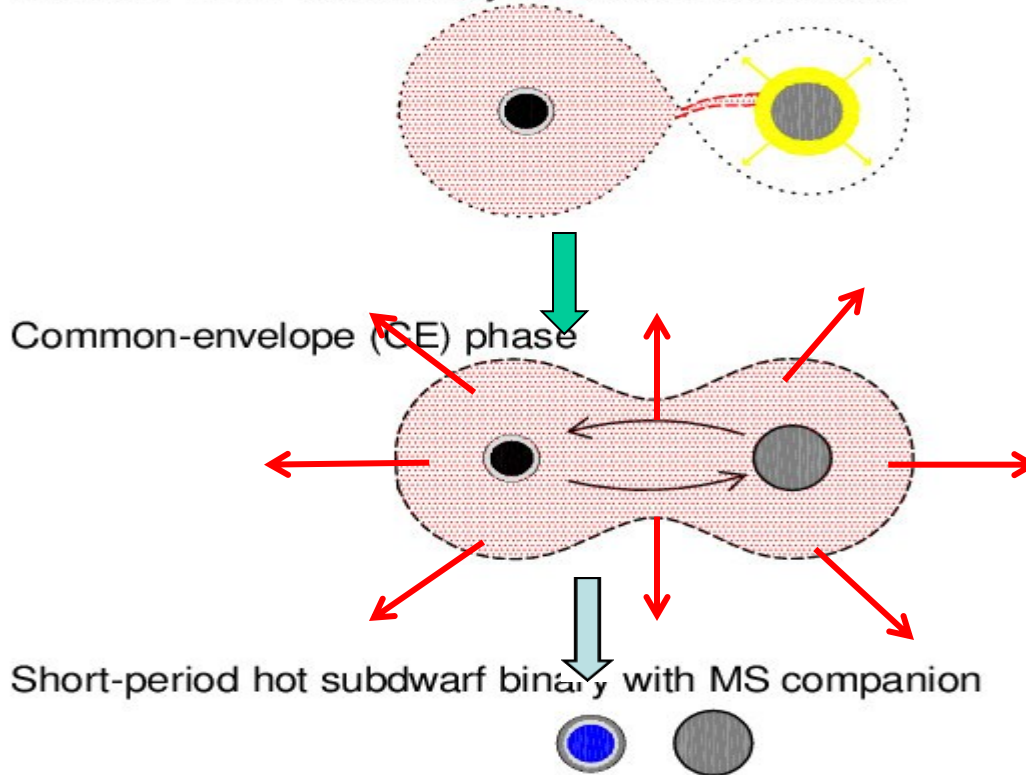


Figure 16. (a) Evolution of mass transfer rate, (b) orbital period as a function of mass, and (c) evolutionary track in the HR diagram, to demonstrate the case of stable RLOF for a binary with a giant donor with an initial mass of $1 M_{\odot}$ and a $0.84 M_{\odot}$ WD companion (Pop I, with overshooting, one-quarter Reimers' wind). The solid curves in (a) and (b) show the evolution before the onset of RLOF, i.e. due to a stellar wind. No stellar wind was included during and after RLOF.

Binary model: CE ejection

Unstable RLOF leads to dynamical mass transfer



Binary model: He WD merger

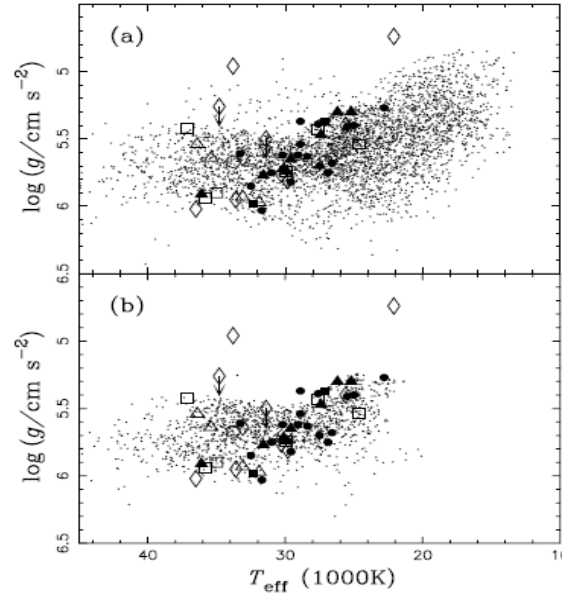
He WD merger (1 or 2 CE phases)



$$M_{\text{sdB}} = 0.40 - 0.65 M_{\odot}$$

The model can reproduce:

- (Log P, Mcomp)
- (Teff, log g)
- P
- Log (g/Teff^4)
- Mass function
- Birthrate
- Number density
- Fraction of sdB+MS
- Fraction of close sdB binaries
-



g diagram for simulation set 2 (the best-fitting model). Dots represent the results of the simulation. Filled circles indicate the position of observed sdB binaries with orbital periods $P_{orb} < 1$ d, solid triangles binaries with $1 < P_{orb} < 10$ d, and solid squares binaries with $P_{orb} > 10$ d. Circles show systems that have radial velocity variations $dV > 10$ km s^{-1} , squares systems with $10 < dV < 20$ km s^{-1} and diamonds systems with $dV < 10$ km s^{-1} , where dV is the maximum difference between radial velocities measured for a particular object. Arrows indicate lower limits for g . (a) does

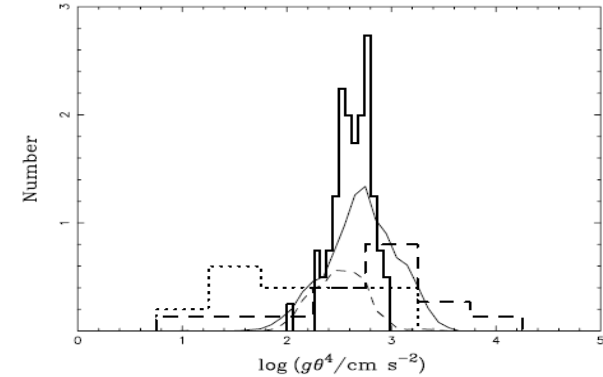


Figure 27. The distribution of $\log(g\theta^4)$, where g is the surface gravity of an sdB star, and $\theta = 5040 \text{ K} / T_{eff}$. The solid histogram is obtained from the 68 sdB stars observed by Saffer et al. (1994), while the dashed and dotted histograms are based on the observations of 15 sdB stars and 10 sdO stars, respectively, by Ulla & Thejll (1998). The solid curve represents the sdB stars from simulation set 2 (the best-fitting model) without consideration of selection effects, while the dashed curve includes the GK selection effect.

Predictions:

Larger mass range

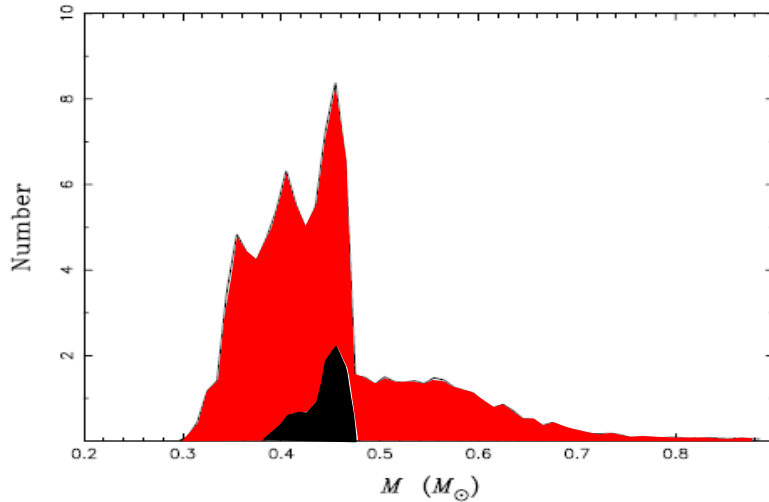


Figure 22. The distribution of masses of sdB stars from all the channels in simulation set 2 (the best-fitting model); solid curve, no selection effects; dashed curve, the GK selection effect; dot-dashed curve, the GK and the strip selection effects; dotted curve, the GK, the strip and the K selection effects.

sdB binaries with long orbital periods

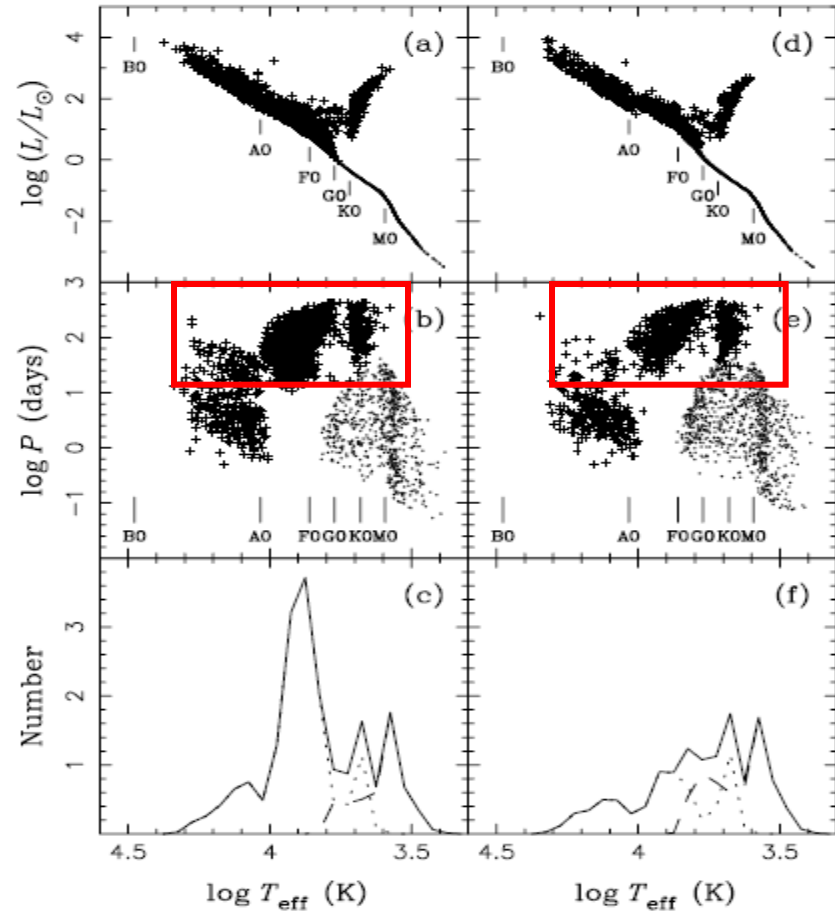


Figure 15. The characteristics of the MS secondaries in sdB binaries from the first CE ejection channel (dots) and the first stable RLOF channel (crosses) for the best-fitting simulations 2 (with $q_{\text{crit}} = 1.5$; left-hand panels) and 8 (with $q_{\text{crit}} = 1.2$; right-hand panels). (a) and (d), HRD with spectral types indicated along the main sequence (based on Zombeck 1990). (b) and (e), orbital period versus effective temperature. (c) and (f), the distribution of $\log T_{\text{eff}}$; solid curves, both channels; dashed curves, the first CE ejection channel only; dotted curves, the first stable RLOF channel.

Outline

- Binary evolution
- Binary-related objects
- **Evolutionary population synthesis**
- Future perspectives

What is EPS?

Evolutionary Population Synthesis (EPS) is an approach to derive the physical parameters of stellar populations of galaxies from the Spectral Energy Distributions (SEDs) observed.



The ingredients of an EPS model:

Star Formation Rate -> how many stars?

Initial Mass Function -> how much mass of each star?

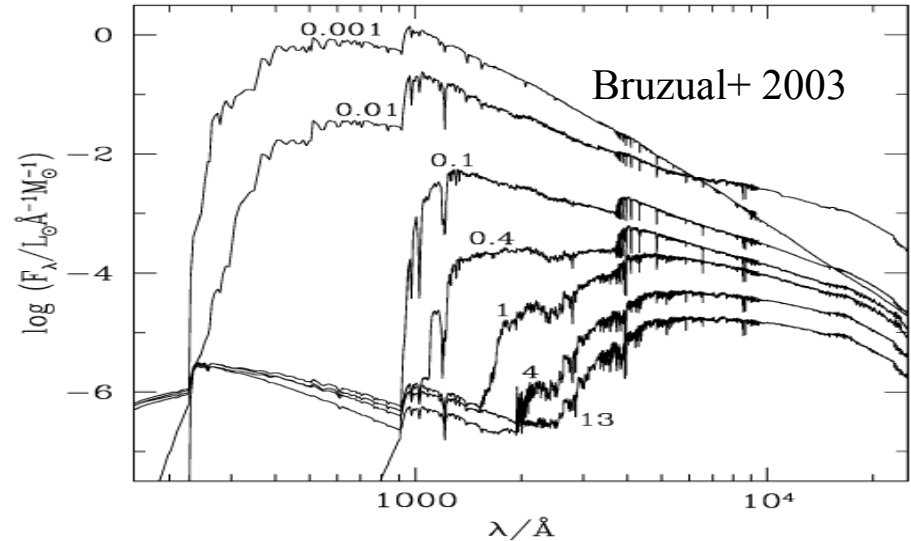
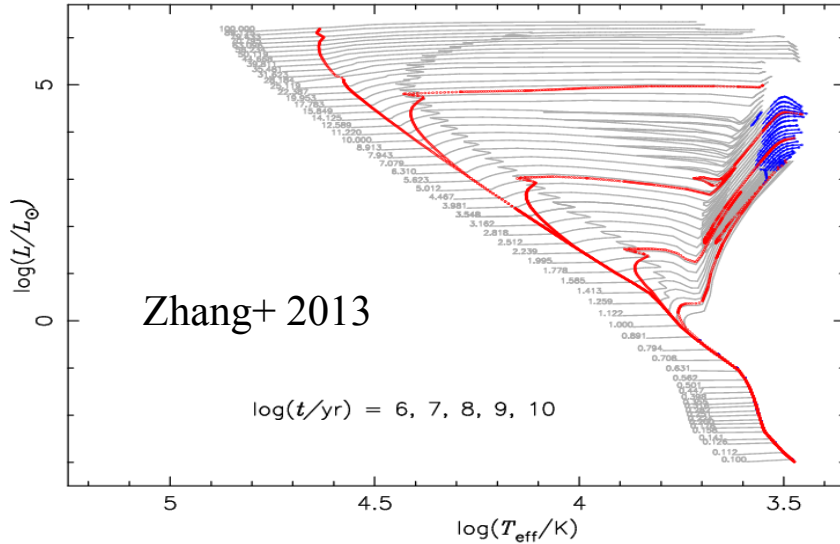
Stellar Evolution Models -> how they evolve?

Model Stellar Atmospheres -> how their spectra?



Integrated SED (the sum of all the spectra of the stars)

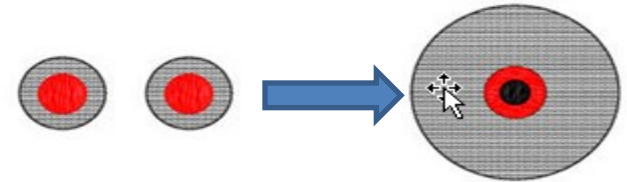
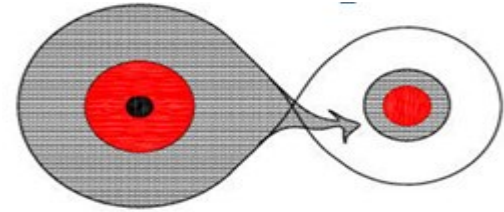
Single Stellar Evolution



Young stellar population has hot spectra, while old population has cold spectra.

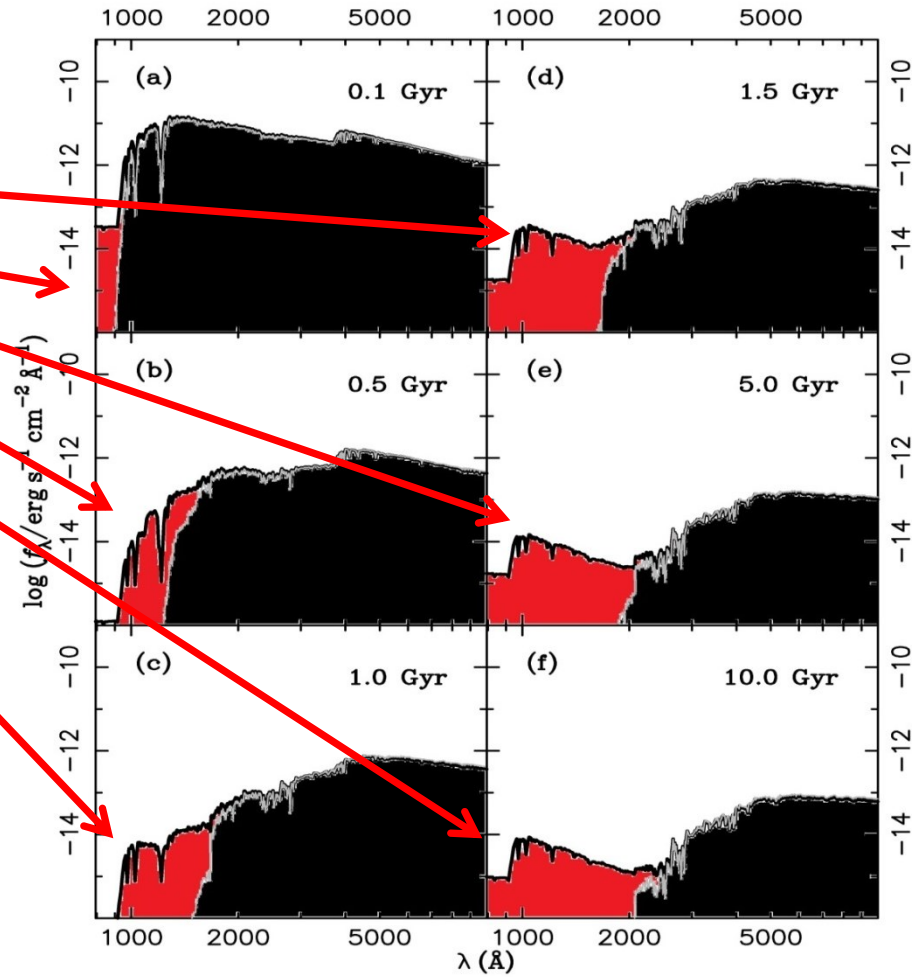
50% are in binaries, how does a binary evolve?

- RLOF removes stellar envelope
→ hot core exposed
→ (core can be ignited)
- A star grows in mass via accretion
→ rejuvenation (hotter)
- Coalescence of a binary
→ a more massive star (hotter)



Binary interactions rejuvenate stars!

Binary
Contribution
to ISED



To derive physical parameters properly from a given SED,

we need binaries !

History of EPS studies

- Tinsley (1968)
- Rapid progress since 90's
- ✓ GISSEL98 or BC03 model (Bruzual+ 1993, 2003)
- ✓ PopSTAR model (Molla+ 2009)
- ✓ STARBURST (Leitherer+ 1999)
- ✓ PEGASE (Fioc+ 1997)
- ✓ **Yunnan model** (Zhang+ 2004, Han+ 2007, Chen+ 2015)



The Yunnan Model

- ✓ stellar evolutionary tracks
(Cambridge stellar evolution code STAR)
- ✓ BaSeL spectral library
- ✓ Binary interactions (blue stragglers, hot subdwarf stars, accreting WDs, rejuvenated stars,

Binaries first included: Zhang+ 2004, A&A, 415, 117

Hot subdwarfs: Han+ 2007, MNRAS, 380, 1098

Blue stragglers: Chen+ 2009, MNRAS, 395, 1822

SFR calibration: Zhang+ 2012, MNRAS, 421, 743

Ionizing sources: Zhang+ 2015, MNRAS, 447, L21

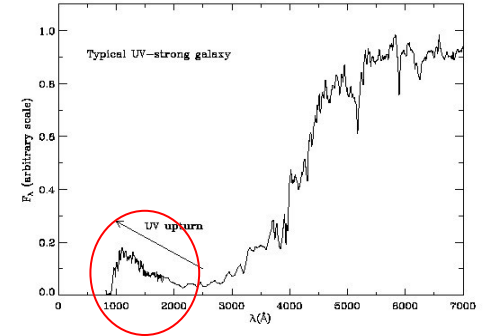
Accreting WDs: Chen+ 2015, MNRAS, 453, 3024

Far-UV Excess of Elliptical Galaxies

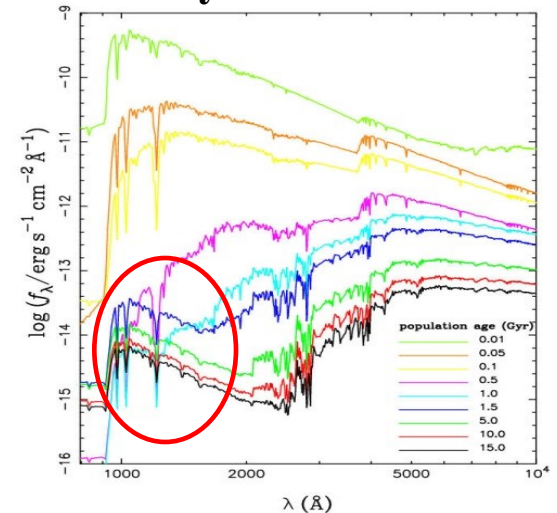
There exists a far-UV excess for elliptical galaxies, known since 1969 but unexplained.

Solving the problem is crucial to our understanding of the structure and evolution of early-type galaxies.

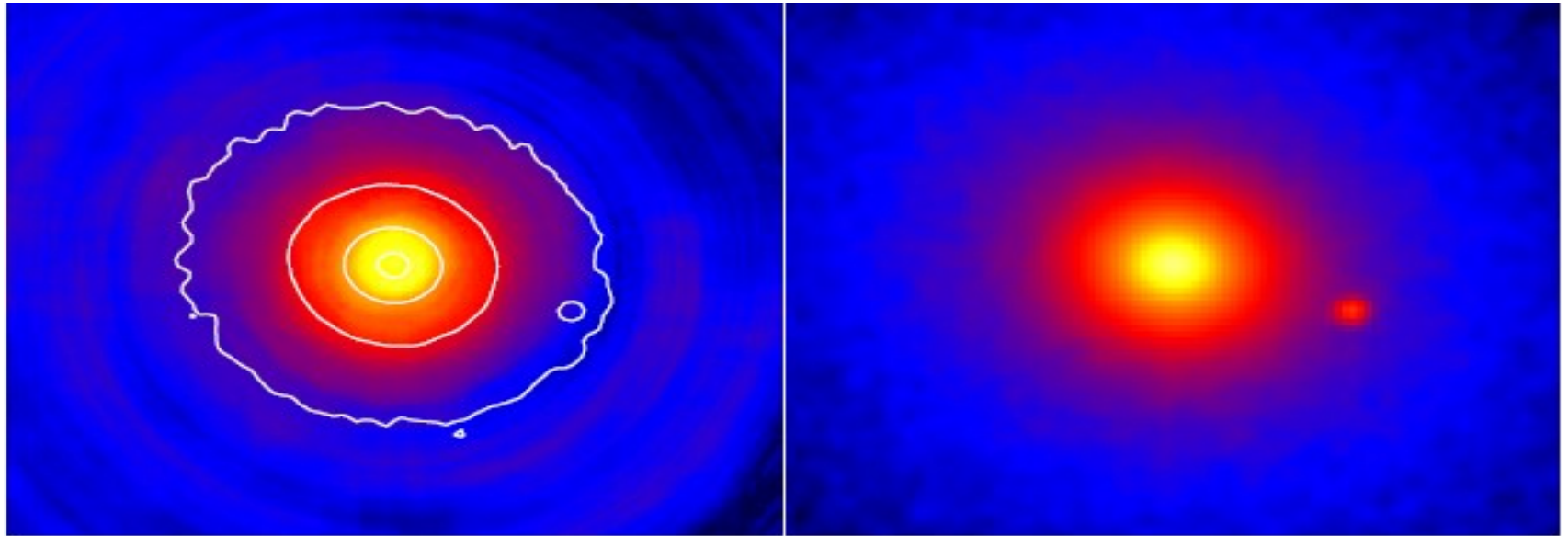
Hot subdwarf stars radiate in far-UV, and we applied our binary model of hot subdwarfs to the study of elliptical galaxies and explained their far-UV light



Binary UV model



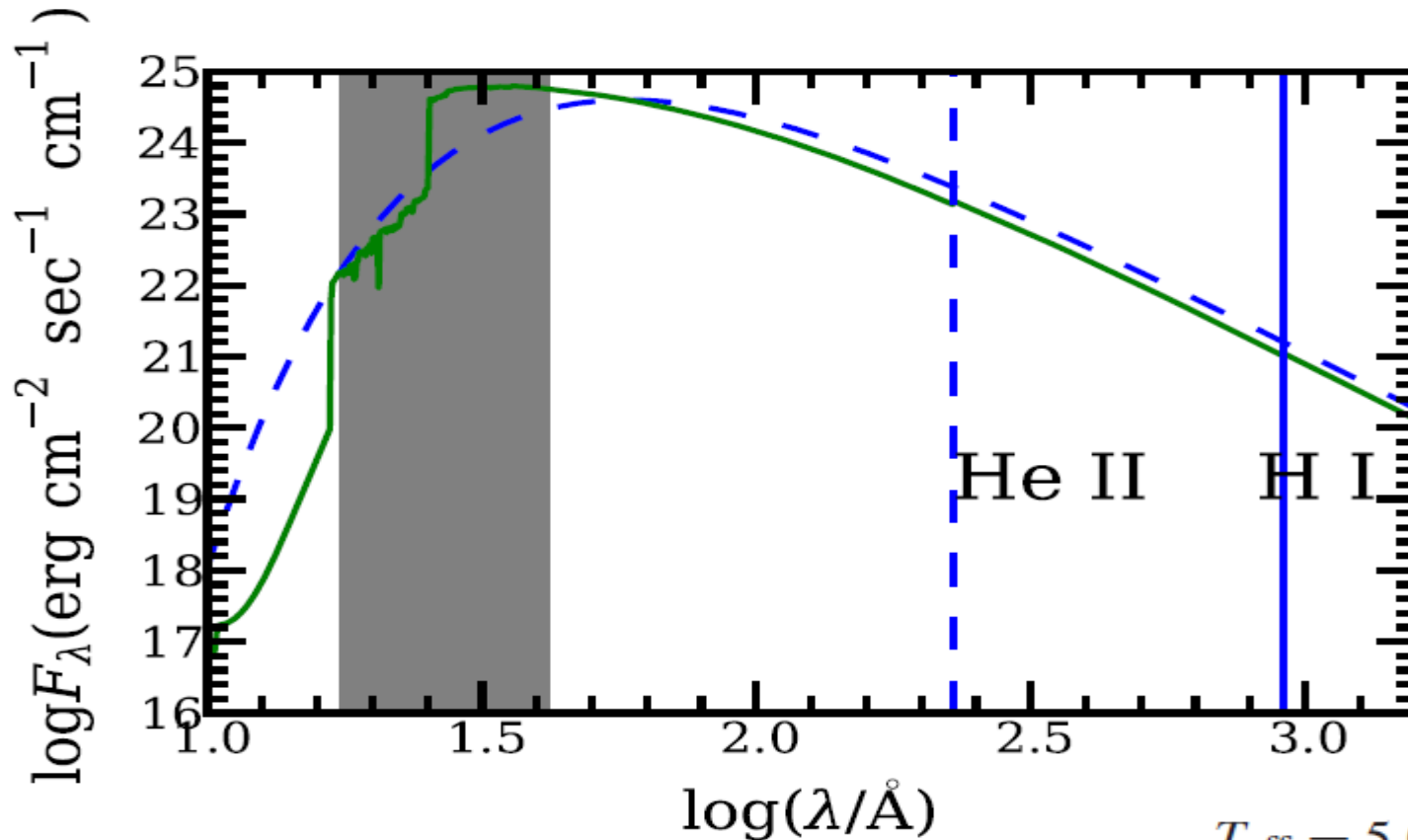
Soft X-ray emission from elliptical galaxies



adaptively smoothed image of the unresolved X-ray emission (detected point source) compared with the near-infrared (*K*-band) image of the galaxy (right panel). Contour lines represent contiguously spaced near-infrared brightness isophotes.

Revnivtsev+ 2007

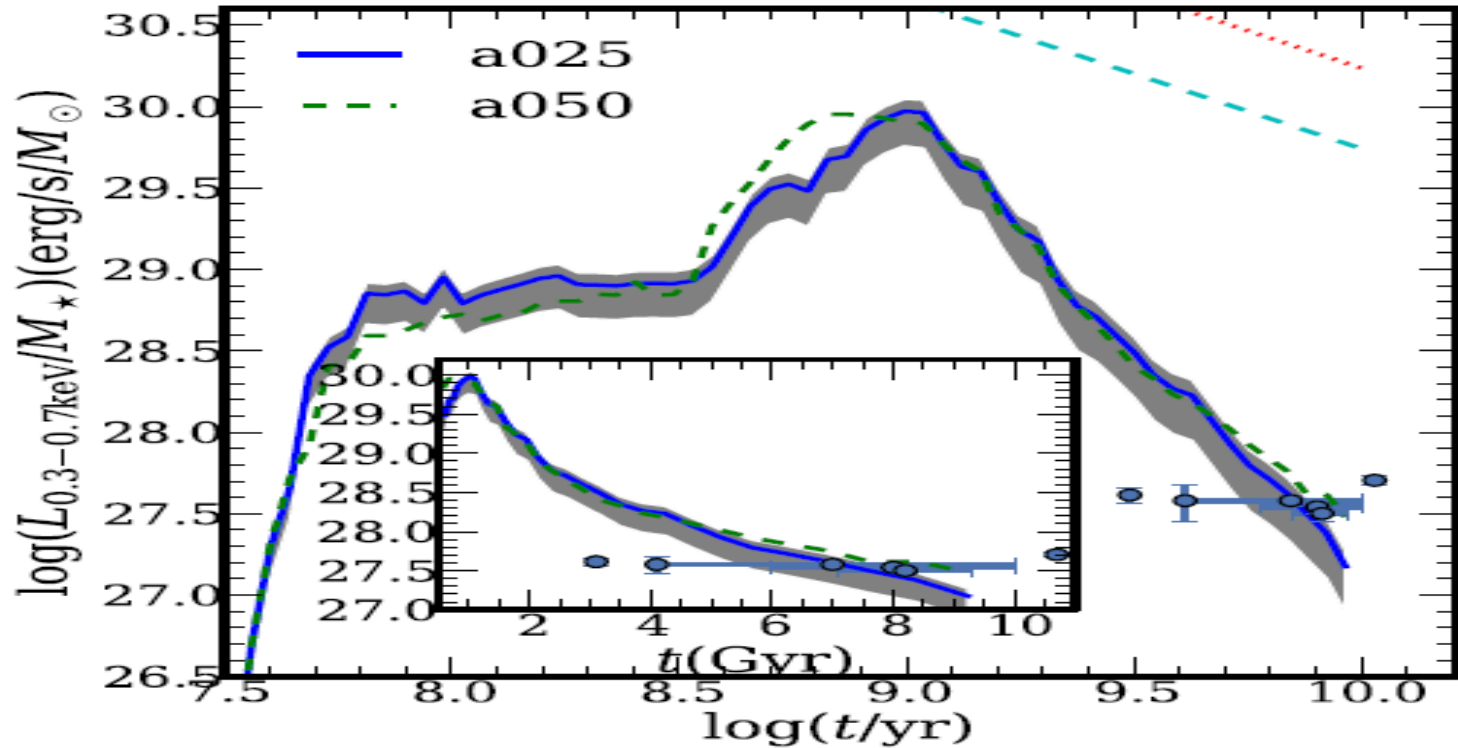
Spectrum of an accreting WD



Chen+ 2014, 2015

$T_{\text{eff}} = 5.0 \times 10^5 \text{ K}$,

Evolution of soft X-ray emission



Chen+ 2014, 2015

Observations: Bogdan+ 2010, Zhang+2012

Yunnan Model:

Stellar mass smaller (by 10%)
fainter (by 0.2mag)
bluer (by 0.2-0.4mag)
Mass-light ratio bigger (by 10-20%)
UV flux larger (by 2 orders of mag)
EUUV flux larger
X-ray emission

Parameters derived:

Age larger (by 20% to a few times)
Metallicity larger (by 20%)
Photometric redshift (much more accurate for blue-faint-galaxies)
Ionizing sources (middle age)
Far-IR bump (?)
SFR (by 60%, 0.2dex)

Zhang+ 2004, 2010, 2012, 2013, 2014
Chen+ 2014, 2015

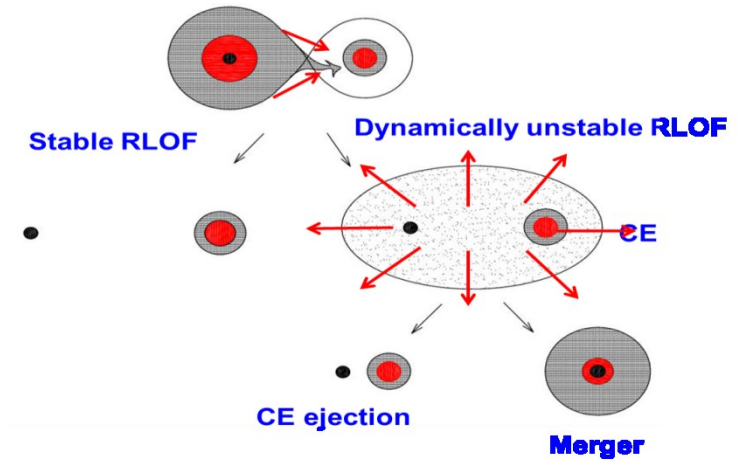
Binary interactions are important for SEDs of old stellar population (e.g. elliptical galaxies) at short wavelength.

Outline

- Binary evolution
- Binary-related objects
- Evolutionary population synthesis
- Future perspectives

(1) Basic problems of binary evolution

- ✓ Ge's adiabatic mass loss model fully established.
- ✓ Symbiotic channel of SNe Ia with Ge's model.
- ✓ Other RG related objects with Ge's model.



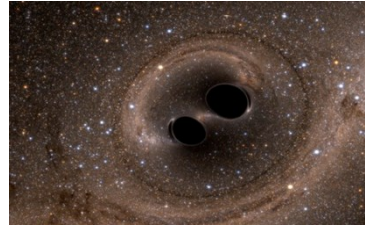
(2) more binary related objects

✓ sdB+NS binaris

✓ BH+BH, NS+NS

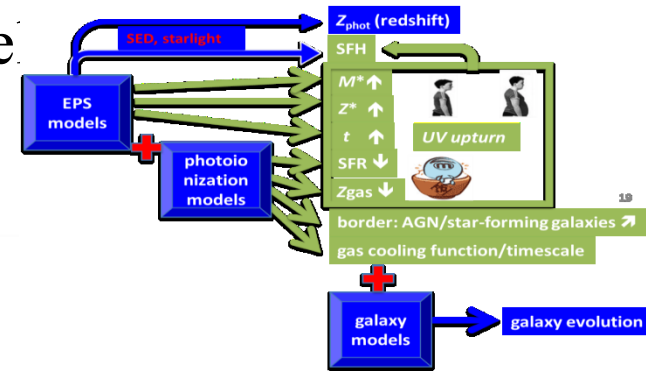
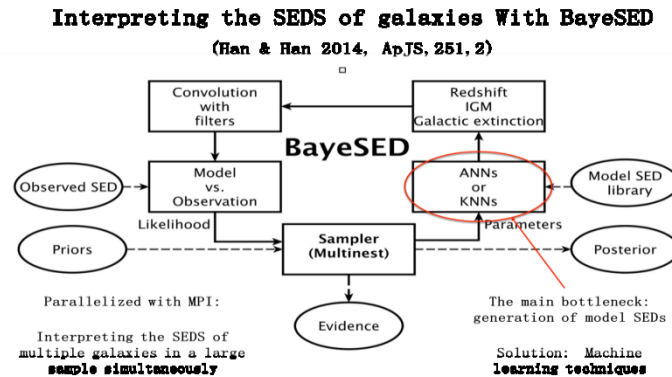
✓ WD+WD

.....



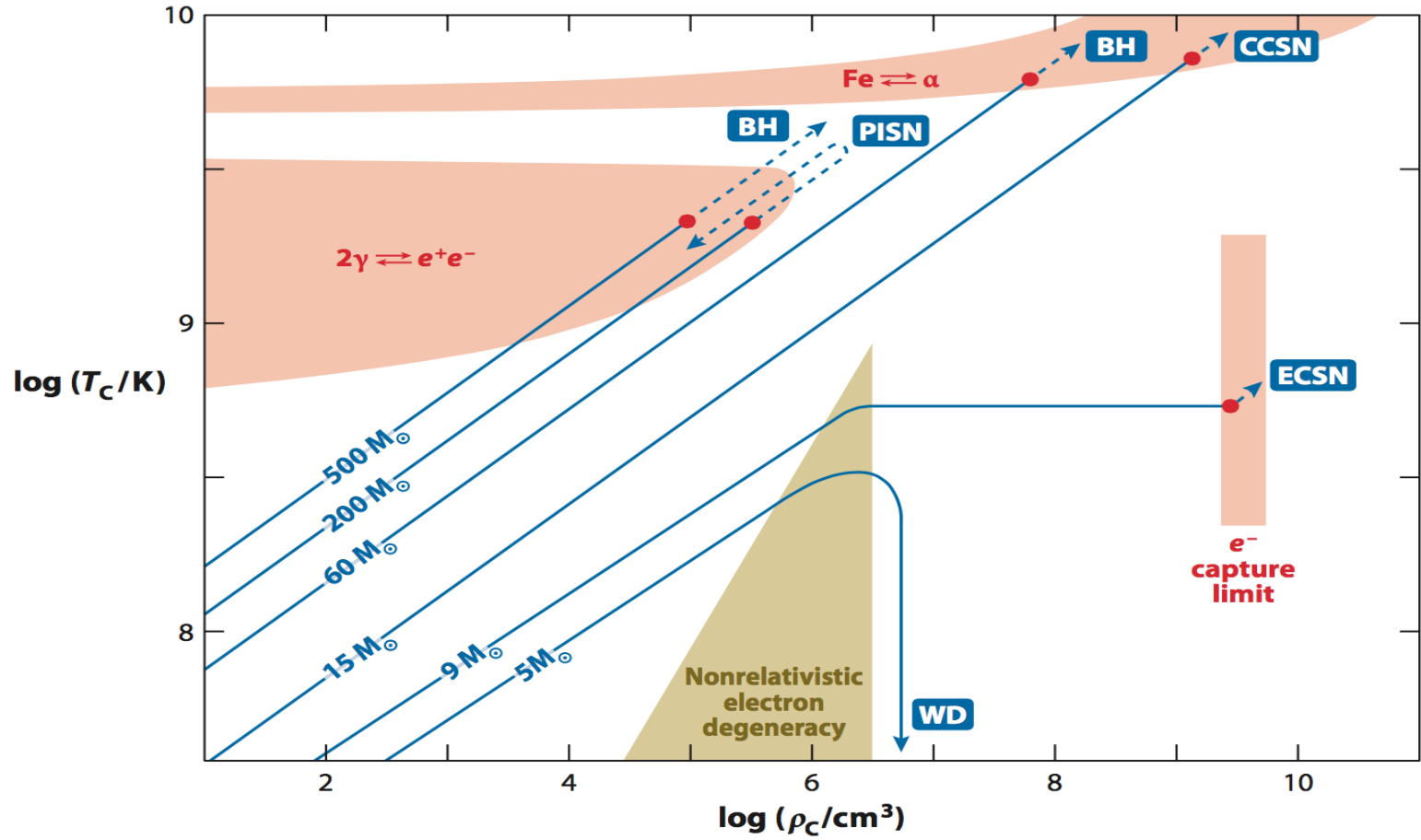
(3) evolutionary population synthesis

- ✓ Extends to EUV, X-ray via accreting WDs
- ✓ +photoionization model + galaxy mode
- ✓ BayeSEI





Thanks !



Science with Inspiral Detections: Maximum BH Mass

1) Maximum BH mass depends on:

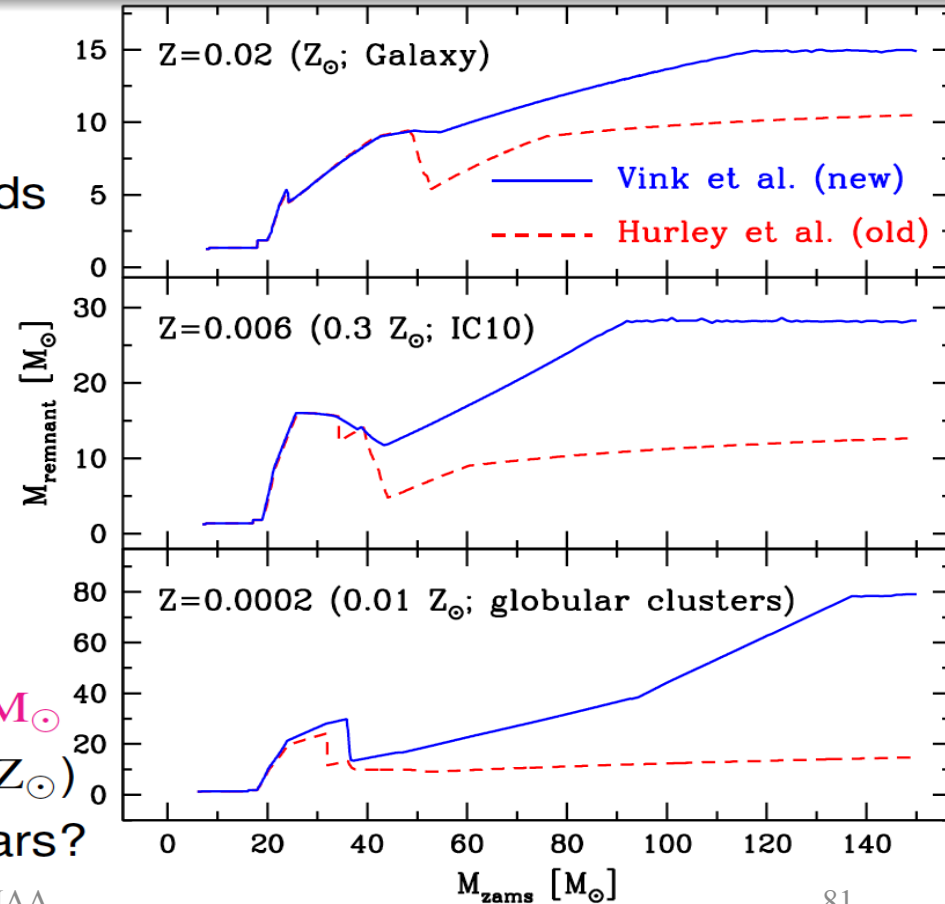
- extent of Initial Mass Function (IMF)
- stellar models and metallicity \rightarrow winds
- SN/core collapse (pair instability)

2) New winds (Vink et al.):

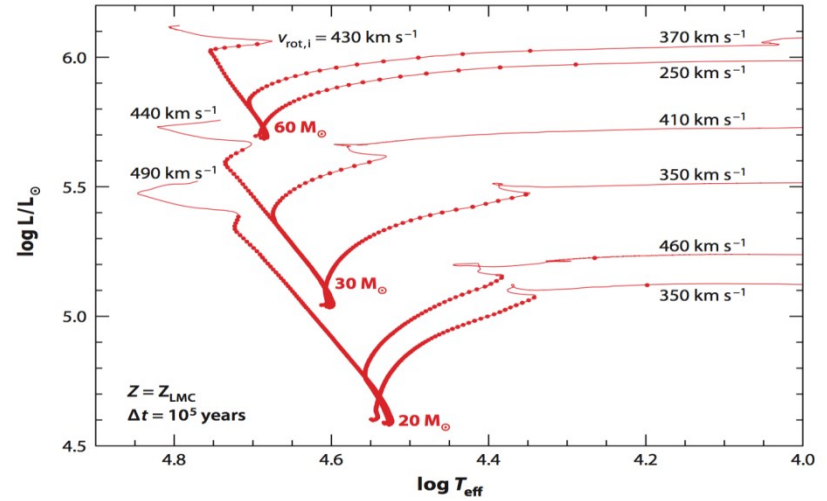
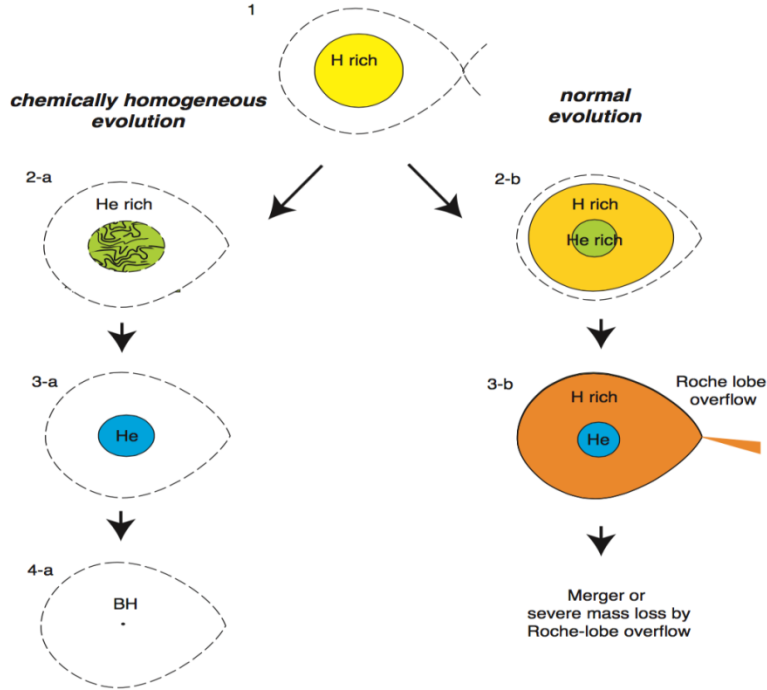
- $Z = 1.0$ Z_{\odot} : max. BH mass: $\sim 15M_{\odot}$
- $Z = 0.3$ Z_{\odot} : max. BH mass: $\sim 30M_{\odot}$
- $Z = 0.01$ Z_{\odot} : max. BH mass: $\sim 80M_{\odot}$

3) Extent of IMF:

- very massive stars found: $200-300 M_{\odot}$
- max. BH mass: $\sim 100-150 M_{\odot}$ ($0.1 Z_{\odot}$)
- but: pair instability disrupts these stars?



Quasi-chemical homogeneous evolution

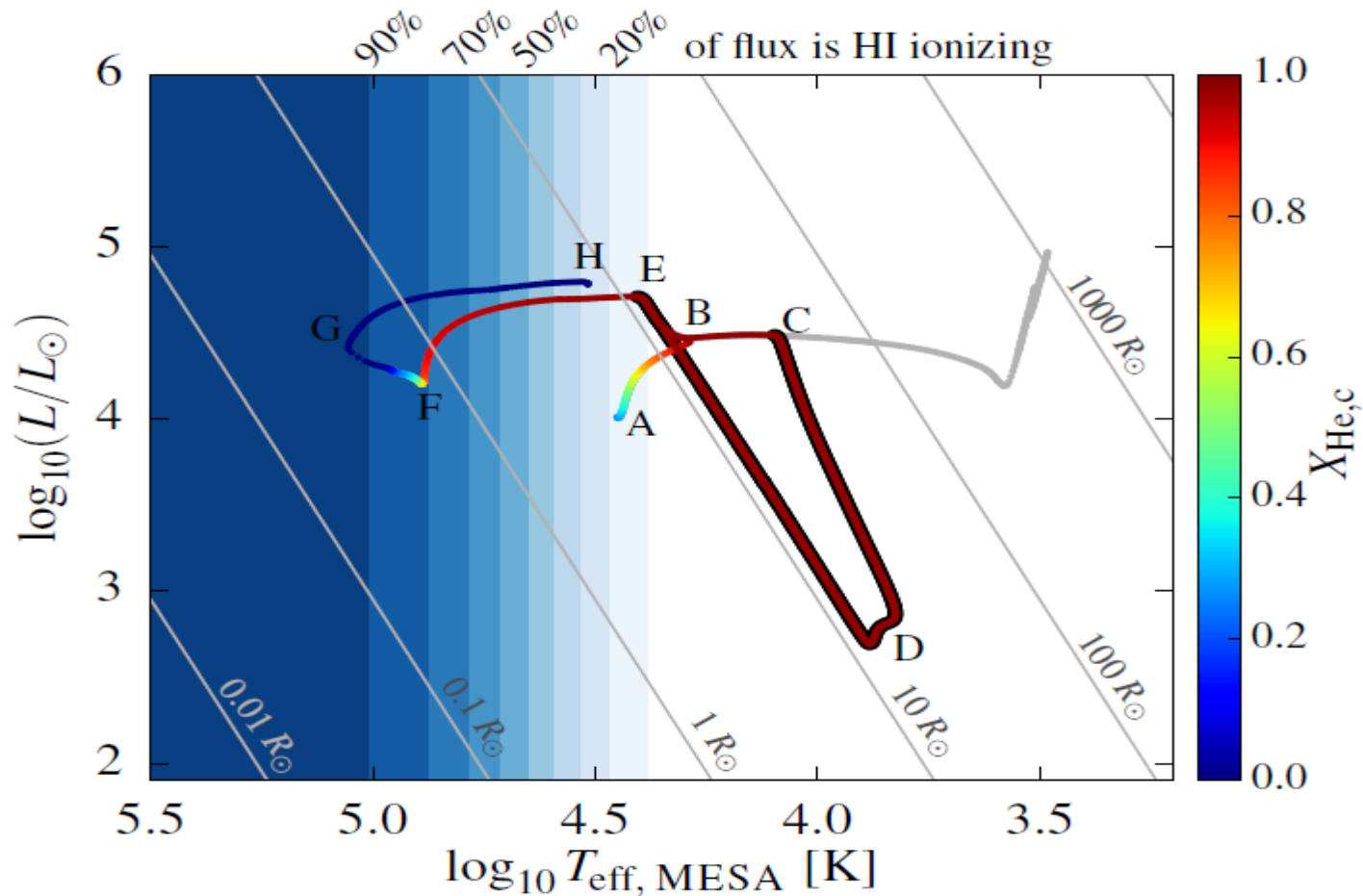


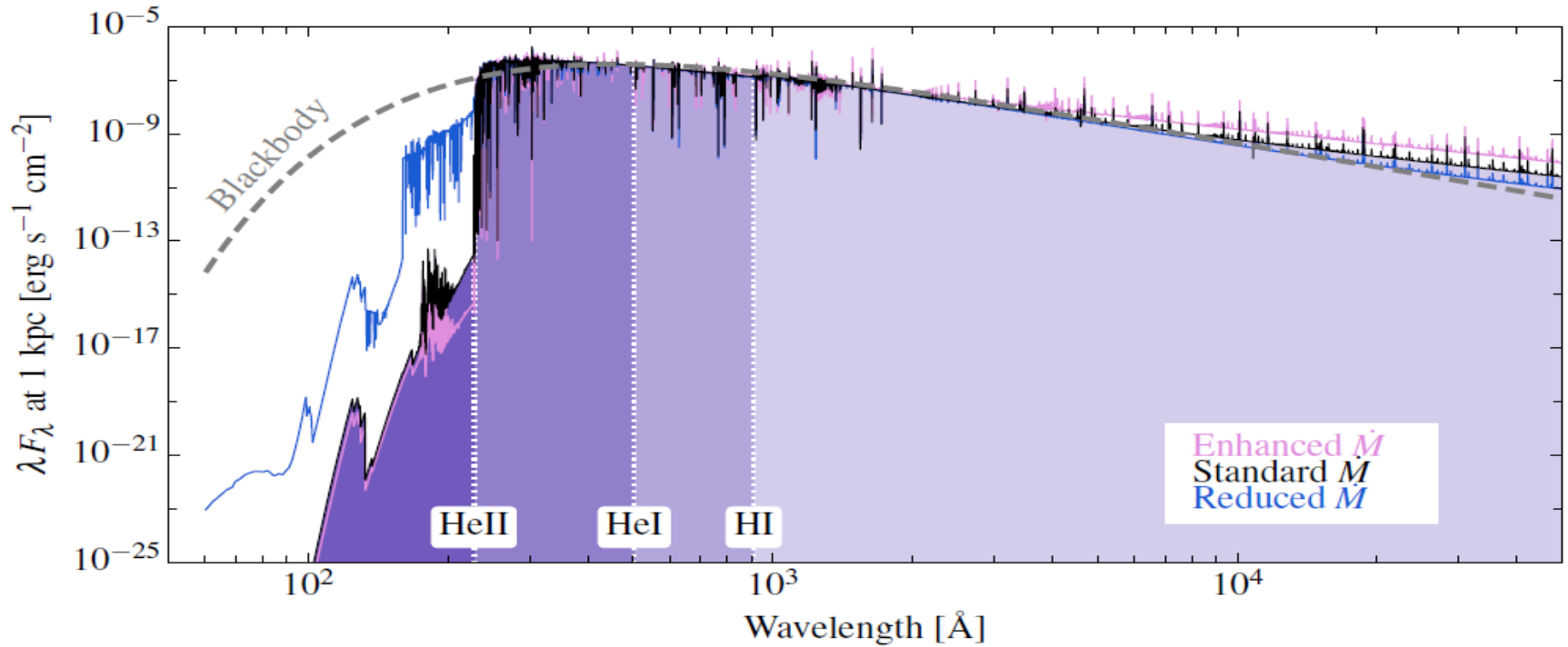
Langer, 2012, ARA&A

Mandel+, 2016, MNRAS

Ionizing spectra of stars that lose their envelope through interaction with a binary companion: role of metallicity

Y. Götberg¹, S. E. de Mink¹, and J. H. Groh²





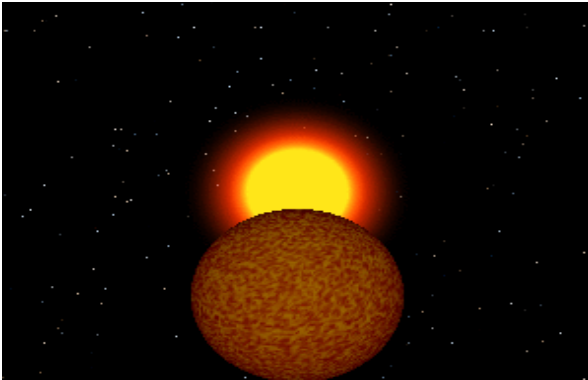
$$\eta = \frac{f_{\text{strip}}}{f_{\text{WR}, 60 M_{\odot}}} \times \frac{Q_{0, \text{strip}}}{Q_{0, \text{WR}}} \times \frac{\Delta t_{\text{strip}}}{\Delta t_{\text{WR}}} \times \left(\frac{M_{\text{strip, init}}}{M_{\text{WR, init}}} \right)^{-2.35}$$

$$= \frac{0.33}{1.0} \times \frac{1.19 \times 10^{48} \text{ s}^{-1}}{2.8 \times 10^{49} \text{ s}^{-1}} \times \frac{1.2 \text{ Myr}}{0.4 \text{ Myr}} \times \left(\frac{12 M_{\odot}}{60 M_{\odot}} \right)^{-2.35} \approx 1.8$$

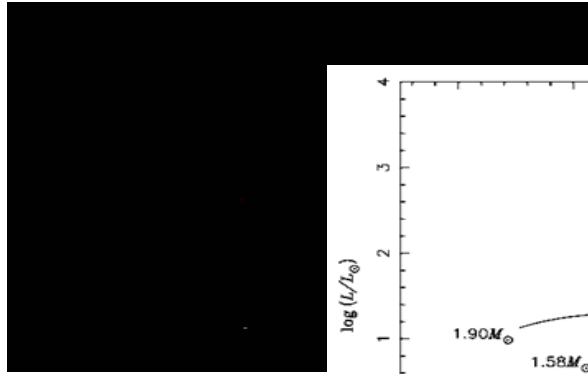
Canonical channels for sdOB formation

(Han+ 2002, 2003)

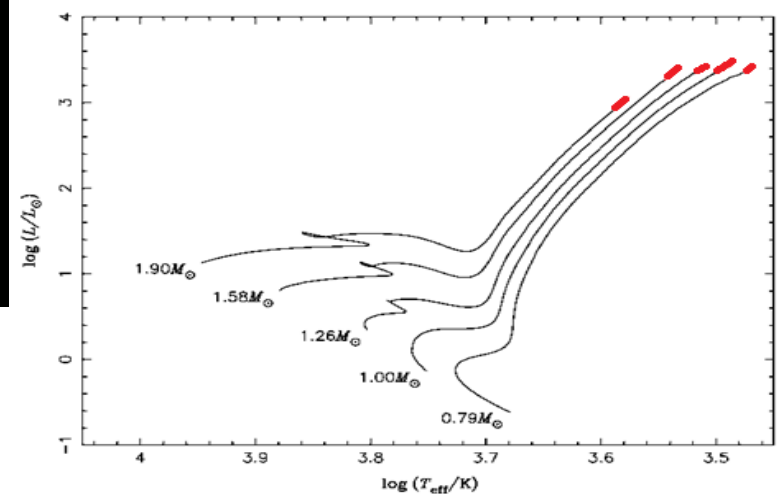
Stable RLOF,



CE ejection,



merger



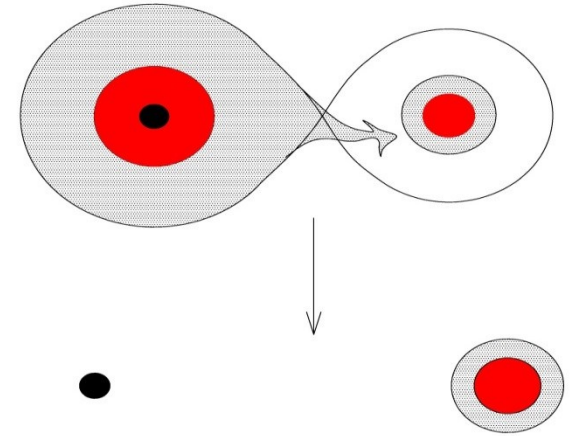
Stable RLOF

During stable RLOF, the donor is Roche-lobe filling.

$M_c \rightarrow R$, $q \rightarrow R_L/A$,
 A can be determined,

The RLOF stops when donor's envelope collapses,
and we then have a q_f - A relation (M_c is fixed)

The relation does NOT depend on dJ



To explain the q - A figure,
 we assume that $M_{1i}=1M_{\text{sun}}$, $R_{\text{FGB}}=0.75\text{AU}$

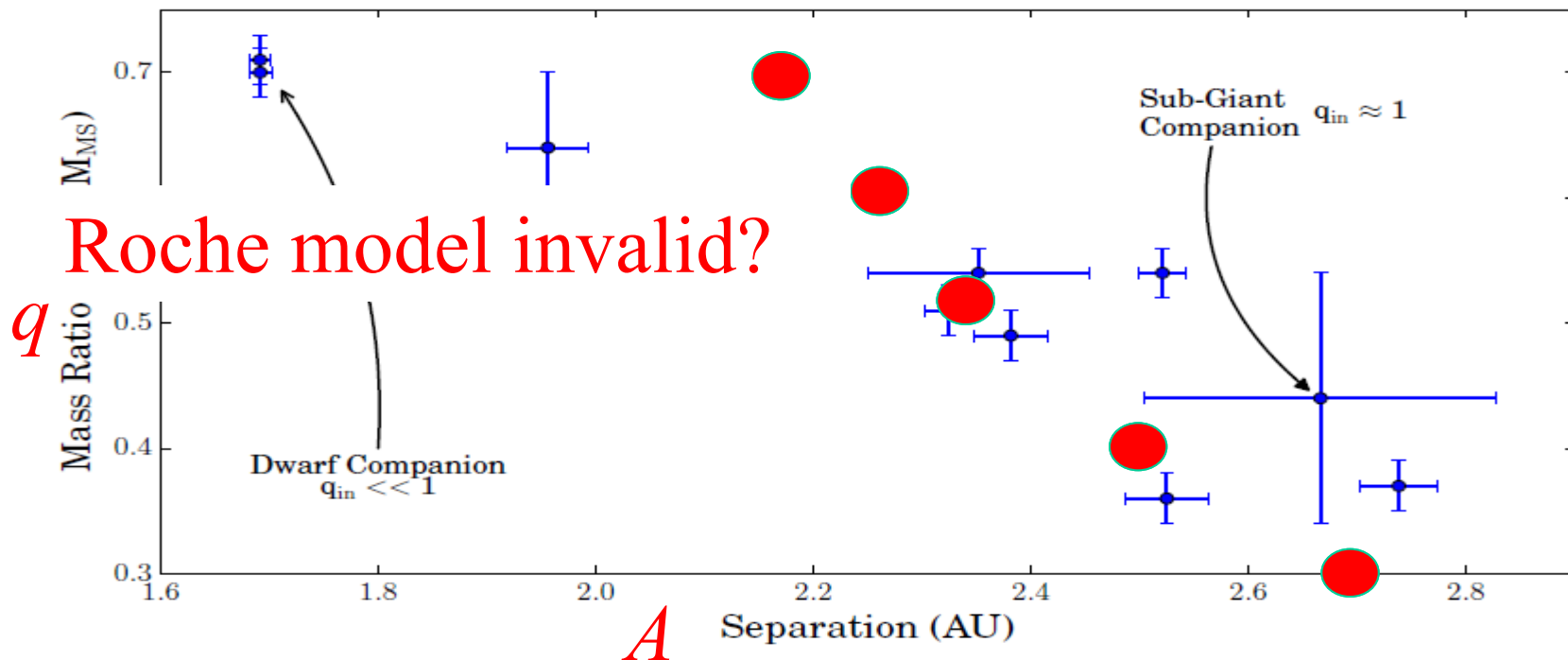
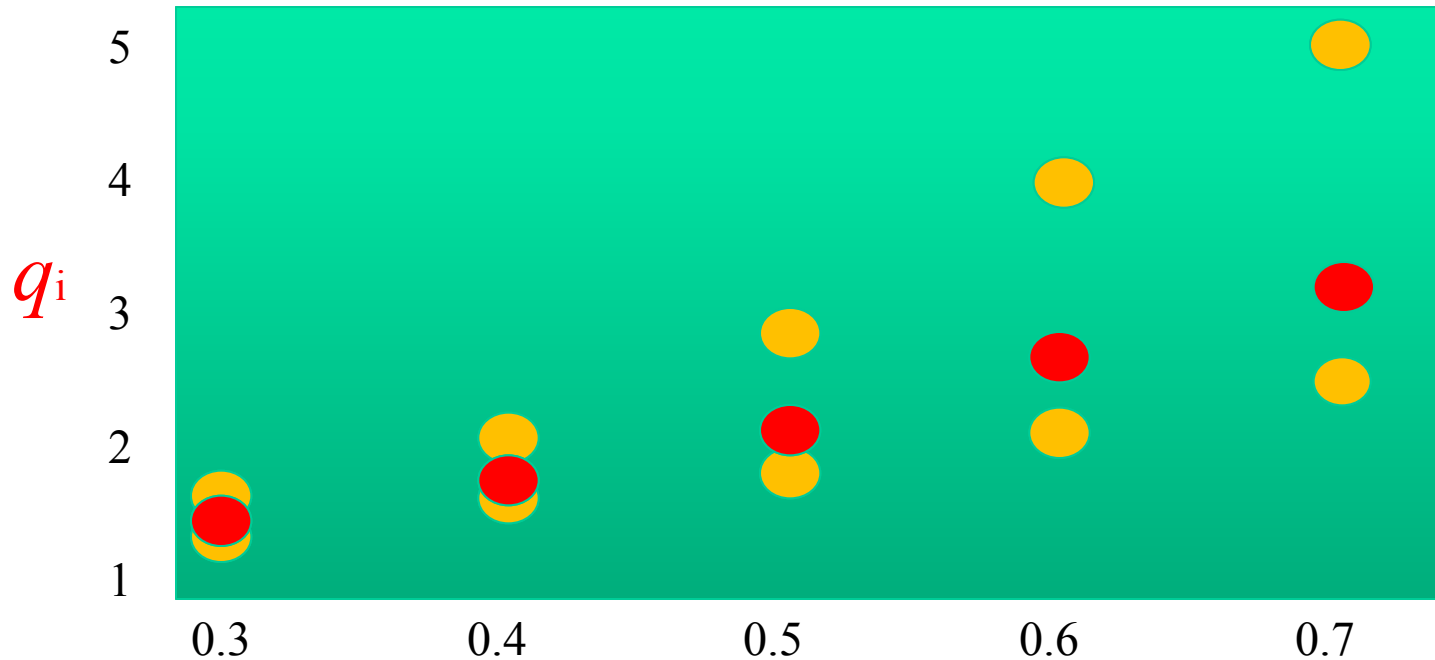


Fig. 2: The separation – mass ratio distribution.

To constrain the stability criterion, q_c ,
we assume that $M_{1i}=1M_{\text{sun}}$, we can derive M_{2i} by
assuming 0%, 50%, 100% of the transferred mass is lost
from the system during RLOF



Hot subdwarf binaries from the MUCHFUSS project

Analysis of 12 new systems and a study of the short-period binary population*

T. Kupfer¹, S. Geier², U. Heber³, R. H. Østensen⁴, B. N. Barlow⁵, P. F. L. Maxted⁶, C. Heuser³,
V. Schaffenroth^{3,7}, and B. T. Gänsicke⁸

unknown companion

dM companion

WD companion

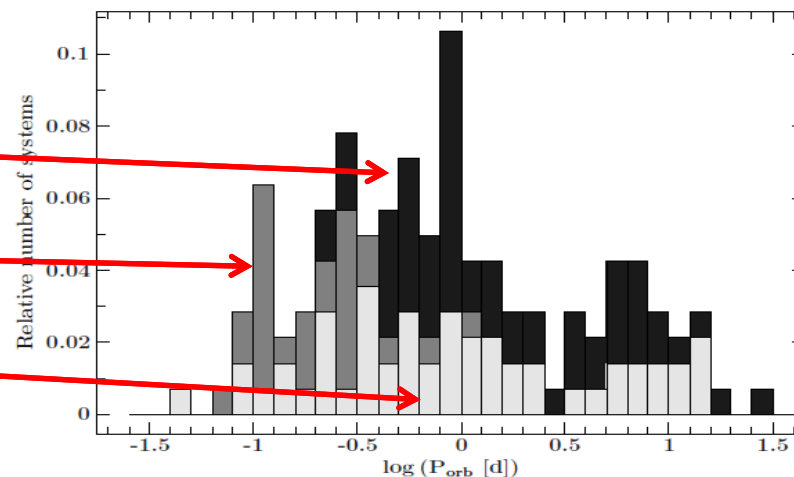


Fig. 7. Period histogram of the full sample. Light grey: WD companions, grey: dM companion, dark grey: unknown companion type.

Chen+, in prep.

